

Competition between nucleon- and $\bar{K}NN$ -cluster correlations in kaonic nuclear systems

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Collaborators:

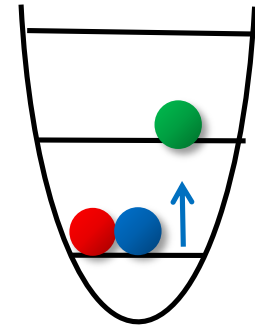
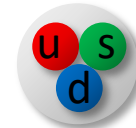
S. Ohnishi, T. Hoshino (Hokkaido → company)

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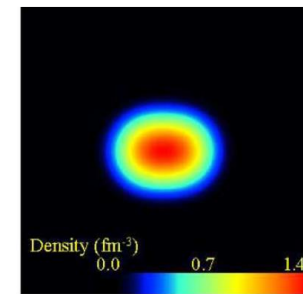
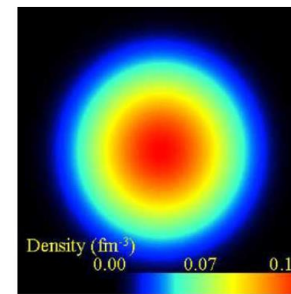
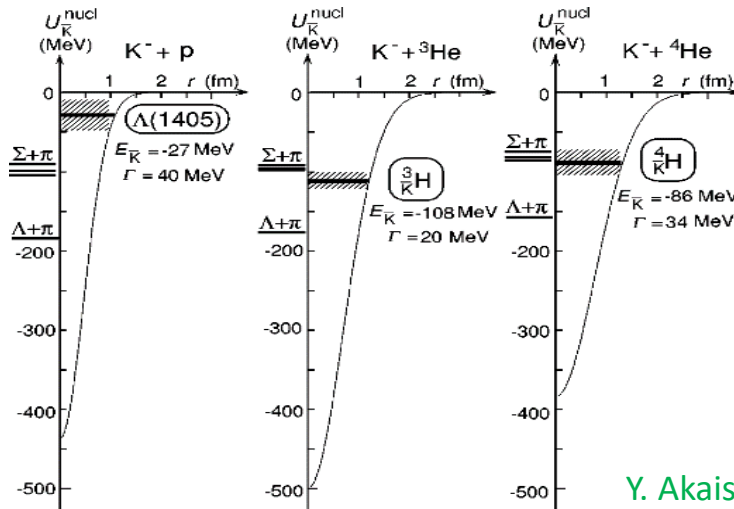
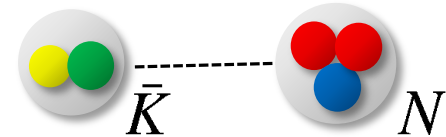
Kaonic nuclei (Nucleus with antikaon)

- $\Lambda(1405)$; $J^\pi=1/2^-$, $S= -1$
 - uds constituent quark model
 - Energy is too high
 - $\bar{K}N$ quasi-bound state
 - **strongly attractive** $\bar{K}N$ interaction



Isgur, Karl, PRD 18, 4187(1978)

Dalitz, Wong, Tajasekaran, PR 153, 1617 (1967)



(a) ${}^3\text{He}$

(b) ${}^3\text{He}K^-$

Dote, et. al., PLB590, 51(2004).

Y. Akaishi, T. Yamazaki, PRC 65, 044005 (2002).

Does Kaonic nucleus really exist? E15 exp. Y. Sada et al., Prog. Theor. Exp. Phys. 2016, 051D01 (2016).

Can such a high density system be produced in laboratory? → $\bar{K}N$ interaction is essential!

Outline

- Precise few-body calculations for kaonic nuclear systems S. Ohnishi, WH, T. Hoshino, K. Miyahara, T. Hyodo, Phys. Rev. C 95, 065202 (2017)
 - Modern $\bar{K}N$ interaction K. Miyahara, T. Hyodo, PRC93 (2016)
 - Structure of kaonic nuclei
 - $\bar{K}NN$ to $\bar{K}NNNNNN$ (7-body)
 - **Clustering**: competition between $4N$ and $\bar{K}NN$ correlations
Variational calculations with correlated Gaussian method

Choice of $\bar{K}N$ interaction

Kyoto $\bar{K}N$ potential

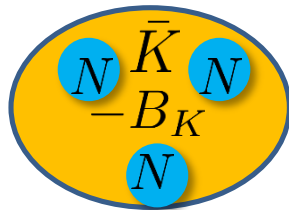
K.Miyahara, T.Hyodo, PRC 93, 015201 (2016)

- **Energy-dependent** $\bar{K}N$ single-channel potential
- Chiral SU(3) dynamics at NLO
- Pole energy: 1424 - 26i and 1381 - 81i MeV Y.Ikeda, T.Hyodo, W.Weise, NPA881 (2012) 98
- **Consistent with the recent kaonic hydrogen data** SIDDHARTA: Bazzi+2012, NPA881

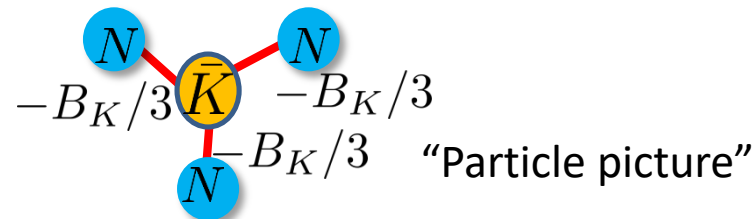
$\bar{K}N$ two-body energy in an N -body system are determined as:

$$\sqrt{s} = m_N + m_{\bar{K}} + \delta\sqrt{s} \quad , \quad -B_K \equiv \langle \Psi | H | \Psi \rangle - \langle \Psi | H_N | \Psi \rangle \quad ,$$

Type I: $\delta\sqrt{s} = -B_K$, Type II: $\delta\sqrt{s} = -B_K / (N - 1)$, for N -body



“Field picture”



A. Dote, T. Hyodo, W. Weise, NPA804, 197 (2008).

Akaishi-Yamazaki (AY) potential

Akaishi, Yamazaki, PRC65, 04400(2002).

- **Energy-independent**
- Reproduce $\Lambda(1405)$ as $\bar{K}N$ quasi-bound state

Kaonic nuclear systems (3 to 7-body)

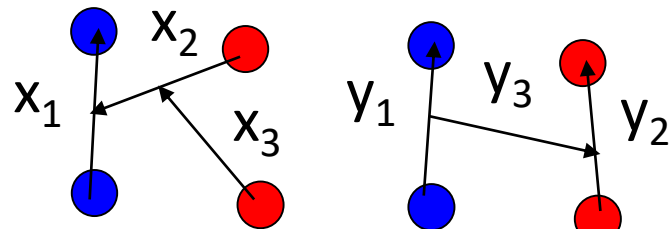
- Hamiltonian

$$H = \sum_{i=1}^{\mathcal{N}} T_i - T_{\text{cm}} + \sum_{i<j}^{\mathcal{N}-1} V_{ij}^{(NN)} + \sum_{i=1}^{\mathcal{N}-1} V_{i,\mathcal{N}}^{(\bar{K}N)} + \sum_{i<j}^{\mathcal{N}} V_{ij}^{\text{Coul.}}$$

- Basis expansion with correlated Gaussian basis:

Formulation for N-particle system

Functional form unchanged under any coordinate trans.



$$\Phi_{SM_S M_T}(\mathbf{x}, \mathbf{A}) = \mathcal{A} \{ \exp(-\tilde{\mathbf{x}} \mathbf{A} \mathbf{x}) \chi_{SM_S} \eta_{M_T} \},$$

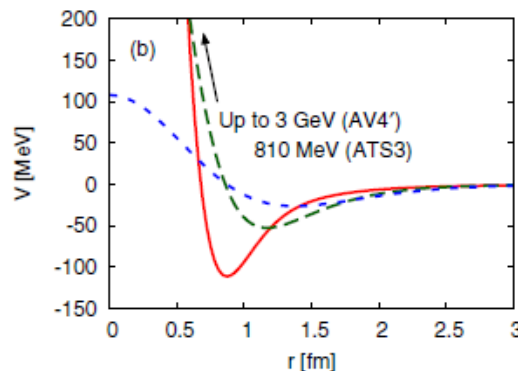
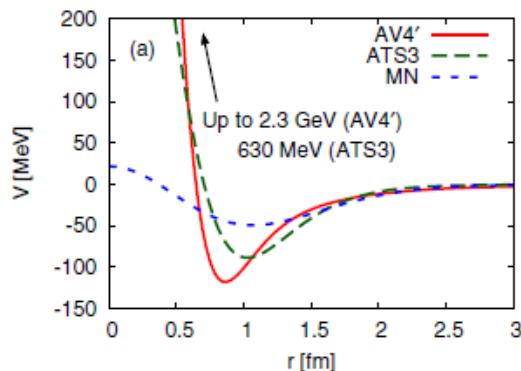
$$\tilde{\mathbf{x}} \mathbf{A} \mathbf{x} = \sum_{i,j=1}^{N-1} A_{ij} \mathbf{x}_i \cdot \mathbf{x}_j$$

$$\mathbf{y} = \mathbf{T} \mathbf{x} \implies \tilde{\mathbf{y}} \mathbf{B} \mathbf{y} = \tilde{\mathbf{x}} \tilde{\mathbf{T}} \mathbf{B} \mathbf{T} \mathbf{x}$$

- Many parameters $\sim (N-1)(N-2)/2 \times (\# \text{ of basis})$

→ Stochastically selected K. Varga and Y. Suzuki, PRC52, 2885 (1995).

- Choice of NN potential (AV4', ATS3, MN)



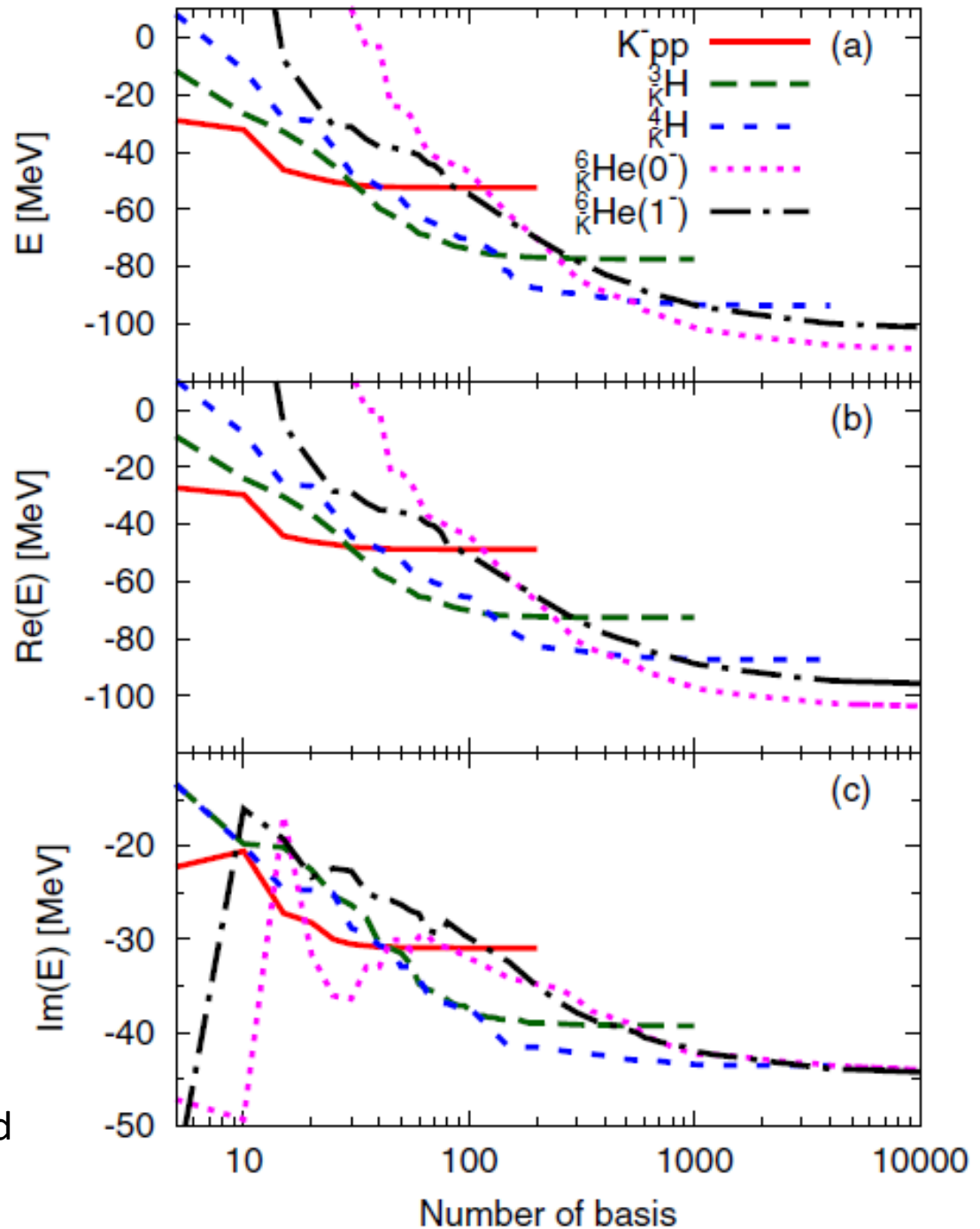
All NN interaction models reproduce the binding energy of s-shell nuclei

Energy curves

- Optimization only with a real part of the $\bar{K}N$ pot.
- Two-body $\bar{K}N$ energy is self-consistently determined
- AV4' NN pot. is employed

Full energy curves

Validity of this approach is confirmed in the three-body (K^-pp) system



Properties of K^-pp (AV4' pot.)

| Model | Kyoto | | AY |
|--|-----------------|-----------------|-------|
| | Type I | Type II | |
| B (MeV) | 27.9 | 26.1 | 48.7 |
| Γ (MeV) | 30.9 | 59.3 | 61.9 |
| $\delta\sqrt{s}$ (MeV) | $-61.0 - i25.0$ | $-30.2 - i23.7$ | |
| P_{K^-} | 0.65 | 0.65 | 0.64 |
| $P_{\bar{K}^0}$ | 0.35 | 0.35 | 0.36 |
| $\sqrt{\langle r_{NN}^2 \rangle}$ (fm) | 2.16 | 2.07 | 1.84 |
| $\sqrt{\langle r_{\bar{K}N}^2 \rangle}$ (fm) | 1.80 | 1.73 | 1.55 |
| $\sqrt{\langle r_N^2 \rangle}$ (fm) | 1.12 | 1.08 | 0.958 |
| $\sqrt{\langle r_{\bar{K}}^2 \rangle}$ (fm) | 1.14 | 1.10 | 0.988 |

Kyoto $\bar{K}N$ pot.

Similar binding energies with
Types I and II $B \sim 26-28$ MeV

$\Gamma \sim 30-60$ MeV

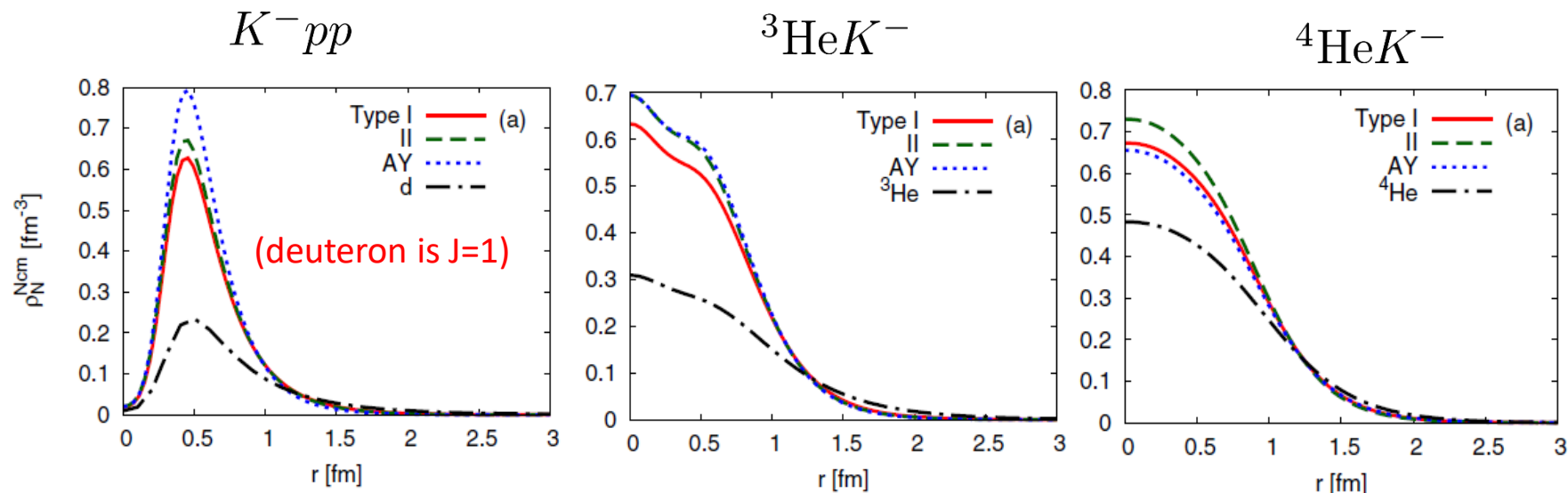
AY pot.

Deeper binding energy ~ 49 MeV
 \rightarrow Smaller rms radii

Systematics of $\bar{K}XN$ ($X=2-6$) systems

| | $\bar{K}NN$ | $\bar{K}NNN$ | $\bar{K}NNNN$ | $\bar{K}NNNNN$ |
|----------------|-------------|--------------|---------------|----------------|
| B (MeV) | 26-28 | 45-50 | 68-76 | 70-81 |
| Γ (MeV) | 31-59 | 26-70 | 28-74 | 24-76 |

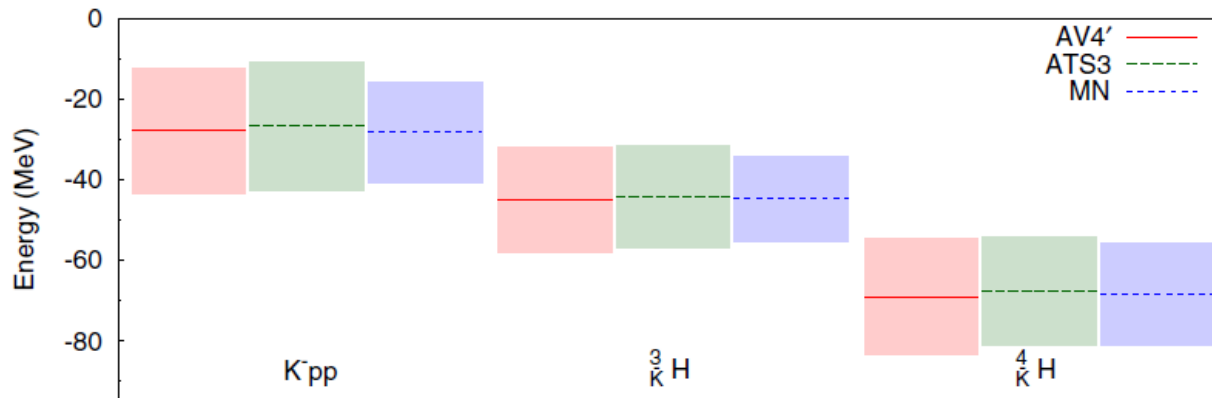
Bound states appear below subthresholds: $B \sim \Gamma$



- Central nucleon density $\rho(0)$ is enhanced by kaon
- Not always proportional to B \rightarrow tail of w.f.
- $\rho(0) \sim 0.7 \text{ fm}^{-3}$ at maximum, ~ 2 times higher than that without \bar{K}
(~ 4 times higher than the saturation dens.)

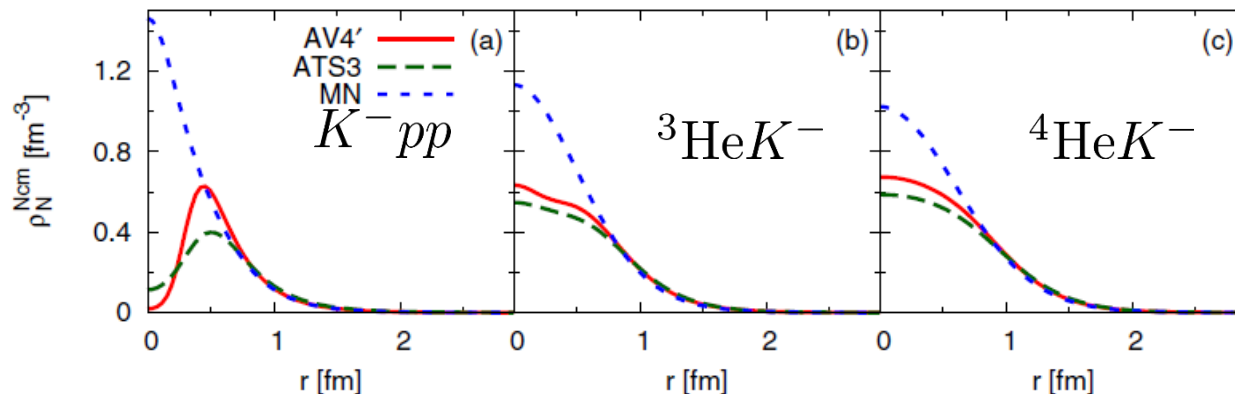
NN interaction dependence

Binding energy and decay width with different NN potential models



Not sensitive to the NN interaction models

Nucleon density distributions

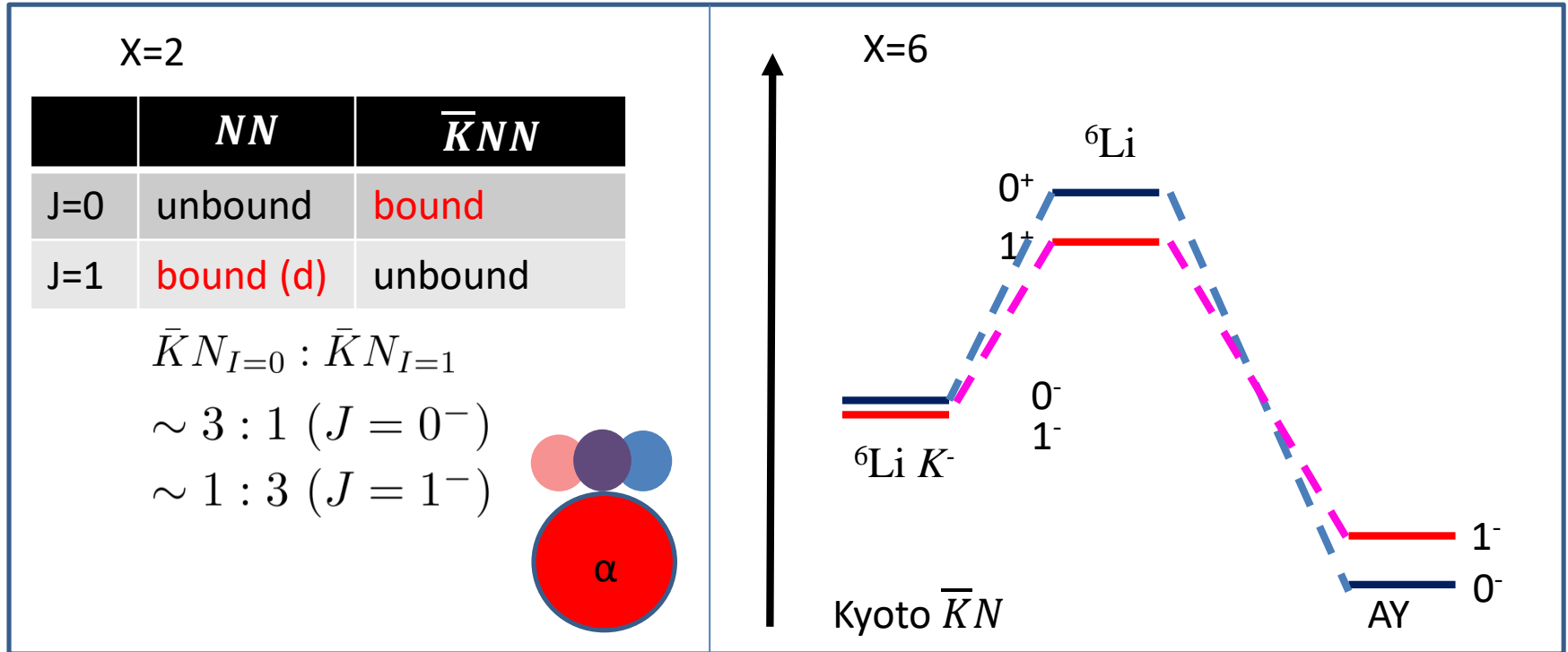


AV4' and ATS3 potential: strong short-range repulsion

MN: weak short-range repulsion

Strong $\bar{K}N$ correlation in $\bar{K}XN$ systems:

Structure of $\bar{K}NN$ & $\bar{K}NNNNNN$



- $\bar{K}N$ interaction in $l=0$ is more attractive than in $l=1$
 - Energy gain in $J=0$ is larger than that in $J=1$
- AY potential in $l=0$ is strongly attractive
 - $J=0$ ground state

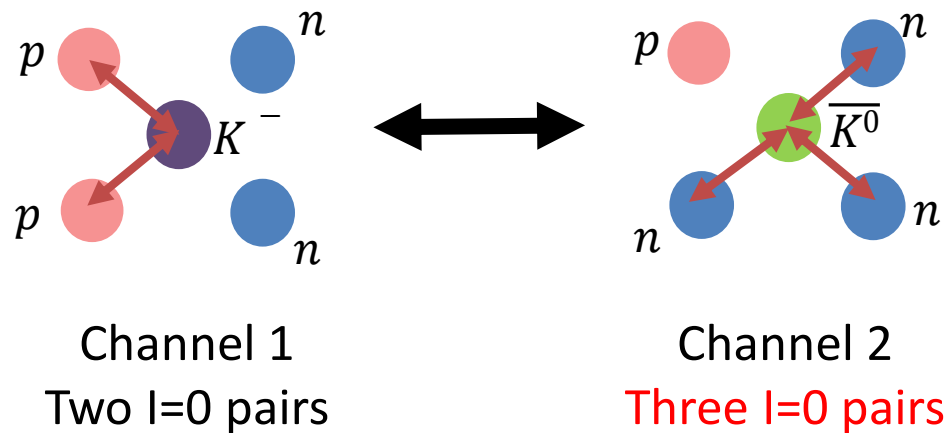
$\bar{K}N$ correlations > NN correlations

$\bar{K}N$ interaction plays a decisive role to determine the structure.

Strong NN correlations: α cluster in $\overline{K}NNNN$

- $J^\pi=0^-, I=1/2, I_z=-1/2$

Mixture of $K^- ppnn$ and $\overline{K}^0 pnnn$ states



- $P_1:P_2=0.93:0.07$
 - Channel 1 can form α cluster
 - $\overline{K}N$ correlations $<$ NN correlations

Summary & perspectives

Precise few-body calculations for kaonic nuclear systems

– Kaonic nucleus (3- to 7-body)

S. Ohnishi, WH, T. Hoshino, K. Miyahara, T. Hyodo,
Phys. Rev. C95, 065202 (2017)

- Central density is increased by ~ 2 times higher (4 times than ρ_0)
 - Soft NN interaction induces too high central densities
 - Possible extensions
 - » Explicit inclusion of the tensor and three-body forces
 - » Explicit coupling of $\pi\Sigma$ and $\pi\Lambda$ channels K. Miyahara, T. Hyodo, W. Weise, arXiv: 1804.08269
- Competition between $\bar{K}NN$ and $4N$ correlations
 - Strong NN correlations in $\bar{K}NNNN$ systems (α correlation)
 - Spin-parity of the g.s. of ${}^6\text{Li } \bar{K}$ (degenerate?, inverted?)

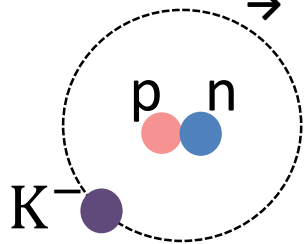
Strength and its isospin dependence of $\bar{K}N$ interaction is essential

→ Kaonic deuterium (3-body) T. Hoshino, S. Ohnishi, WH, T. Hyodo, W. Weise, Phys. Rev. C 96, 045204 (2017)

- Prediction of the energy shift of the kaonic deuterium

$$\Delta E - i\frac{\Gamma}{2} = (670 - i508) \text{ eV},$$

Promising observable to constraint the $l=1$ component of the $\bar{K}N$ interaction
complimentary to the kaonic hydrogen data (SIDDHARTA: Bazzi+2012, NPA881)



Kaonic deuterium