

Determining the ANCs for $^{13}\text{C} \rightarrow ^{12}\text{C} + n$ for Nuclear Astrophysical Studies

Joel Greenberg, Princeton University
REU at the TAMU Cyclotron Institute

Background

The Big Picture

- The gamma ray signature from the decay of ^{22}Na (in nova explosions) has not been observed by gamma-ray telescopes
- One possible explanation is that its precursor, ^{22}Mg , may be depleted by radiative proton capture ($^{22}\text{Mg}(p,\gamma)^{23}\text{Al}$)
- The experiment aims to determine the ANC \rightarrow the astrophysical S-factor and, thus, an accurate estimate of how much ^{22}Na should be expected

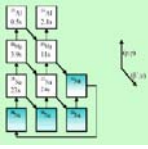


Figure 1. Section of the chart of nuclides representing nuclear reactions and β -decays related to the process ^{22}Na cycle. The stable nuclei in the solid circles have been studied.

In a Direct Radiative Capture Process

$$\sigma \propto |M|^2 \quad M = \langle \Phi^{(13}\text{Al}) | \hat{O} | \Phi^{(22}\text{Mg}) \chi^{(n)} \rangle$$

$$= \langle \psi_{\text{Mg},p}^{(13}\text{Al}) | \hat{O} | \chi^{(n)}(kr) \rangle$$

In a Direct Transfer Reaction

$$\left(\frac{d\sigma}{d\Omega} \right) \propto |M'|^2$$

Where the Transition Amplitude is

$$M' = \langle \Phi^{(12}\text{C}) \Phi^{(23}\text{Ne}) \chi^{(n)} | V_{\text{ex}} | \Phi^{(13}\text{C}) \Phi^{(22}\text{Ne}) \chi^{(n')} \rangle$$

Integrating over the internal coordinates

$$M' = \langle I_{\text{Mg}}^{(22}\text{Ne}, \chi_f^{(n)}) | V_{\text{ex}} | I_{\text{C}}^{(12}\text{C}, \chi_i^{(n')}) \rangle$$

Furthermore, if the transfer is peripheral

$$I_{\text{Mg}}^{(22}\text{Ne}, \chi_f^{(n)}) \approx C_A^{(n)} \frac{W(2kr_M)}{r_M}$$

Thus, for the overall reaction

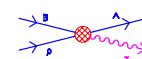
$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{exp}} = (C_{\text{ex}}^{(13}\text{Al}, ^{13}\text{C})})^2 (C_{\text{ex}}^{(22}\text{Ne}, ^{22}\text{Ne})})^2 \left(\frac{d\sigma/d\Omega}{(b_{\text{ex}}^{(13}\text{Al})})^2 (b_{\text{ex}}^{(22}\text{Ne})})^2} \right)$$

The $C_{\text{ex}}^{(13}\text{Al}, ^{13}\text{C})}$ ANC can be easily determined in a separate neutron exchange reaction $^{13}\text{C}(^{12}\text{C}, ^{13}\text{C})^{12}\text{C}$ where the system vertices are identical before and after the reaction

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{exp}} = (C_{\text{ex}}^{(13}\text{Al}, ^{13}\text{C})})^2 \left(\frac{d\sigma/d\Omega}{(b_{\text{ex}}^{(13}\text{Al})})^2} \right)$$

Using Mirror Systems

- Normally a ^{22}Mg beam would be used and the proton transfer rate would be measured directly
- Instead, we
 - Use a ^{22}Ne beam and measure the neutron transfer rate
 - This indirectly measures the original system's mirror system, and uses charge symmetry to back out the appropriate parameters to relate the two



$$\text{Mirror System Relationship}$$

$$^{22}\text{Al} \rightarrow ^{22}\text{Mg} + p$$

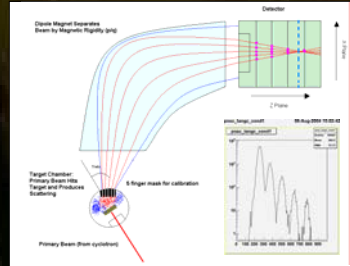
vs.

$$^{22}\text{Ne} \rightarrow ^{22}\text{Ne} + n$$

Experimental Setup

The experiment

- K500 superconducting cyclotron produces ^{22}Ne beam at 264 MeV and bombards $100 \mu\text{g}/\text{cm}^2$ ^{13}C target
 - In order to determine the $^{12}\text{C}(^{13}\text{C}, ^{13}\text{C})^{12}\text{C}$ ANCs, a ^{12}C beam is produced at 127 MeV and bombards an identical target
- The scattered particles then enter the MDM (multipole-dipole-multipole) Spectrometer



The Detector



- The remaining particles enter the Oxford detector which consists of:
 - 4 resistive wires to measure the x position and angle of the incoming particles
 - An ionization chamber filled with isobutane to measure the ΔE (energy loss) in the chamber



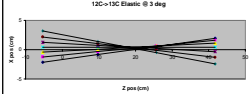
- A plastic scintillator at the back to measure the residual energy

Data Analysis

Data analysis step 1:

Calibration

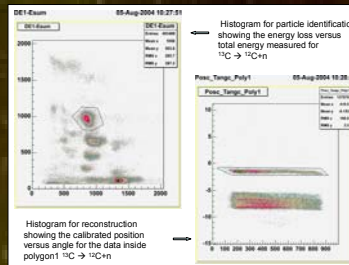
- A ^{197}Au target is first placed in the target chamber for calibrating the position of the wires in the detector
- A mask is placed at the exit of the target chamber so that the particles entering are confined to a 5×0.1 deg slits 0.77 deg apart
- The resulting experimental data can then be used in conjunction with a Raytracing code to determine a proper position calibration for the wires as well as determine the focal plane of the beam
- A similar setup is used with the elastic scattering on the real target to get an angular calibration for the detector



Data analysis step 2:

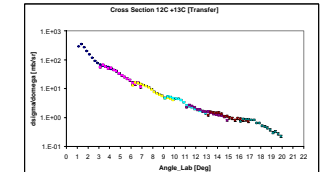
Getting Good Data

- The calibrations are used to produce histograms of the calibrated data
- Events from many reaction products are visible in the data and may overlap
- Within the code, conditions are set in order to select only the desired data
- Multiple histograms and graphs are used for purposes of particle identification and reconstruction



Determining the Experimental Cross section

- The raw data results in an angular distribution which is then normalized to give an absolute cross section

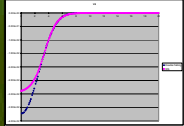


Theoretical Calculations

Calculating the Theoretical Cross Section

- Computer programs (JLM, Ptolemy...) are used to calculate:
 - the nucleon density for the incoming and outgoing channels
 - the double-folded interaction potential
- The double-folded potential is then fit with a Woods-Saxon potential by varying the geometric parameters as well as the normalization constants**

**The results of elastic scattering are to be used to determine the optical model parameters



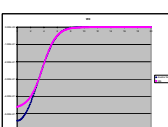
A Woods-Saxon potential is then fit to the double-folded potential (in the region of interest) allowing for the extraction of the optical model parameters

$$U(r) = -V_0 f(r, R, a) - iW_0 f(r, R, a)$$

$$\text{where } f(r) = \frac{1}{1 + e^{\frac{r-R}{a}}}$$

A program using the Hartree-Fock method is used to calculate the nucleon densities and a double-folded potential is calculated

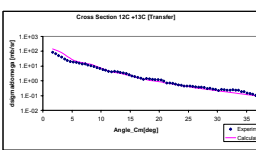
$$U(r) = \iint \rho_1(r_1) \rho_2(r_2) v_{12}(r_1 - r_2) dr_1 dr_2$$



Conclusions

Putting it All Together

- With the Woods-Saxon optical model parameters, a theoretical cross section is then calculated and normalized by comparing it to the experimental results



ANCs

- The proportionality constant between the experimental and calculated cross sections is known as the Spectroscopic Factor

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{exp}} = S_b S_c \left(\frac{d\sigma}{d\Omega} \right)_{\text{DBRBA}} \Rightarrow \left(\frac{d\sigma}{d\Omega} \right)_{\text{exp}} = S^2 \left(\frac{d\sigma}{d\Omega} \right)_{\text{DBRBA}}$$

- The Spectroscopic Factor is directly related to the system ANC

$$S = \frac{C^2}{b^2} \quad \text{Where } C = \text{system ANC}$$

$$b = \text{single particle ANC}$$

→ The preliminary result for the $^{13}\text{C} \rightarrow ^{12}\text{C} + n$ system ANC is: $C^2 \approx 2.1 \text{fm}^{-1*}$

*Value found in literature = 2.55 fm⁻¹