

A 2° Level of ${}^8\text{B}$ and the ${}^7\text{Be}(p, \gamma){}^8\text{B}$ S -factor

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In a recent ${}^7\text{Be} + p$ elastic scattering experiment, Gol'dberg *et al.* [1] found evidence for a wide s -wave 1° or 2° level of ${}^8\text{B}$ at an excitation energy of $E_x = 3.0$ MeV. They speculate on how such a level would affect the low-energy S factor for the ${}^7\text{Be}(p, \gamma){}^8\text{B}$ reaction, which plays a crucial role in the solar neutrino problem. Measurements of the S factor at laboratory energies ($E_{c.m.} > 100$ keV), where $E_{c.m.}$ is the relative p - ${}^7\text{Be}$ energy, have to be extrapolated to the effective energy in the sun's interior ($E_{c.m.} = 20$ keV) using a calculated energy dependence, and this could be changed if a fit to the properties of this ${}^8\text{B}$ level were imposed.

Gol'dberg [2] has identified the level as 2° with a width between 1 and 4 MeV, and has given the experimental energy resolution as 0.25 MeV. A 2° level near 3 MeV may not be unexpected; for example, van Hees and Glaudemans [3], in their shell-model calculations, predict a 2° state in ${}^8\text{Li}$, and in its mirror nucleus ${}^8\text{B}$, at 2.91 MeV (unless otherwise indicated, energies are excitation energies). Most direct measurements of the ${}^7\text{Be}(p, \gamma){}^8\text{B}$ S factor have been confined to the region with $E_x < 2$ MeV.

We here, using R-matrix formulae with parameter values chosen to ${}^7\text{Li} + n$ data, consider the role of 2° states of ${}^8\text{B}$ in the calculated S factors for the ${}^7\text{Be}(p, \gamma){}^8\text{B}$ reaction, and the modification of the S factor caused by the lowest 2° state to be at an excitation energy of 3 MeV. In the R-matrix calculation [4], the spectroscopic factor for the lowest 2° state of ${}^8\text{Li}$ or ${}^8\text{B}$ was taken from a shell-model calculation,

and the energy of the state was then determined by fitting the ${}^7\text{Li} + n$ scattering length $a_2 = -3.63$ fm. Now that an experimental value of the resonance, $E_r = 3.0$ MeV [1], is available for the energy of the state in ${}^8\text{B}$, it seems more reasonable to make use of that, giving the eigenvalue $E_1 = 3.0$ MeV, and to adjust the spectroscopic factor in order to fit a_2 . The resultant value of the resonance width Γ^0 lies within Gol'dberg's suggested range of 1 – 4 MeV [2] for channel radius $a = 4.0$ fm, and is somewhat larger for the larger values of a . Attempts to fit the data with a much smaller value of Γ^0 for the 2° level of ${}^8\text{B}$ at 3 MeV (say 2 MeV or less) would immediately have difficulty in fitting the ${}^7\text{Li} + n$ scattering length a_2 . Apart from that, the 1^+ and 2^+ spectroscopic factors for ${}^8\text{Li}$ and the thermal-neutron capture data could be fitted, but the predicted $F_{n\gamma}$ at higher energies would be higher than the experimental values, and the predicted ${}^7\text{Be}(p, \gamma){}^8\text{B}$ $E1$ S factor would become larger than the experimental values near 3 MeV. The calculated total S_{17} -factor and the s -wave 2° S factor containing the contribution from the resonance at 3 MeV are shown in Fig. 1. The low-energy S factor is shown in Fig. 2 in comparison with the most-recent direct measurements by Hammache *et al.* [5] and Hass *et al.* [6]. Fig. 2 also shows the earlier direct measurements by Vaughn *et al.* [7] and Filippone *et al.* [8]. Recent experimental values of the $E1$ S factor determined from the Coulomb dissociation of ${}^8\text{B}$ [9,10] tend to lie higher and are also shown in Fig. 2. Recapping we can conclude that the width of the 2° level at 3 MeV

in ${}^8\text{B}$ needs to be near the upper end of the range 1 – 4 MeV suggested by Gol'dberg [2], in order to fit all the data. Assumption of a width as small as, say, 2 MeV makes it impossible to fit the ${}^7\text{Li} + n$ data and leads to unreasonably large values at higher energies for the ${}^7\text{Be}(p,\gamma){}^8\text{B}$ S factor. A measurement of the width of the 2^- level from a reaction such as ${}^7\text{Be}({}^3\text{He},d){}^8\text{B}$ [11] seems desirable. More accurate new direct measurements of the S_{17} factor at energies above 2 MeV seem also important.

References

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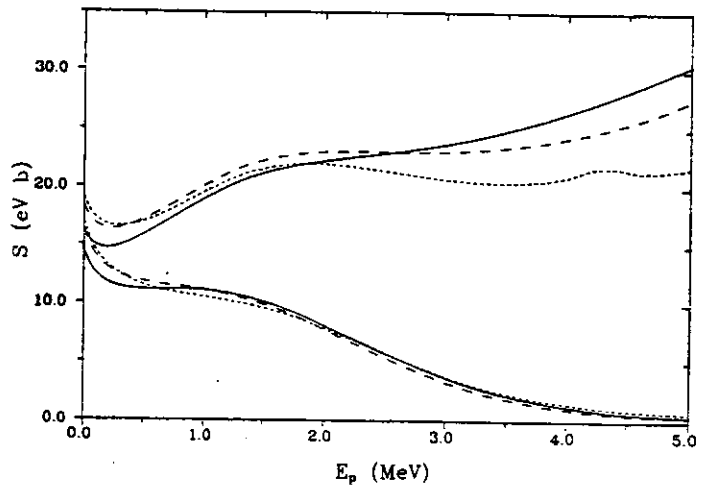


Figure 1. Calculated ${}^7\text{Be}(p,\gamma){}^8\text{B}$ $E1$, S factor as a function of proton energy, for channel radius $a = 4.0$ fm (solid curve), 5.0 fm (dashed curve) and 6.0 fm (short dashed curve), in the four-level, four-channel R-matrix approximation. The s -wave 2^- contribution is shown in the lower curves.

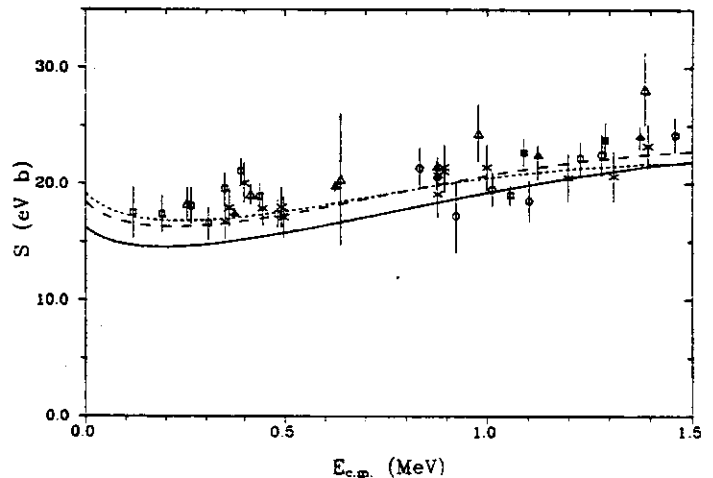


Figure 2. ${}^7\text{Be}(p,\gamma){}^8\text{B}$ S factor as a function of ${}^7\text{Be} + p$ c.m. energy. The experimental points are from Hammache *et al.* [5] (crosses), Hass *et al.* [6] (solid squares), Vaughn *et al.* [7] (circles), Filippone *et al.* [8] (open squares), Kikuchi *et al.* [9] (closed triangles), and Iwasa *et al.* [10] (open triangles). The curves are the predicted $E1$ S factor for $a=4.0$ m (solid curve), 5.0 fm (dashed curve) and 6.0 fm (dotted curve), in the one-level, one-channel R-matrix approximation.