



IRENA



Photoneutron Calibration Standard: extracting cross sections of GDR response in the heavy nucleus ^{169}Tm

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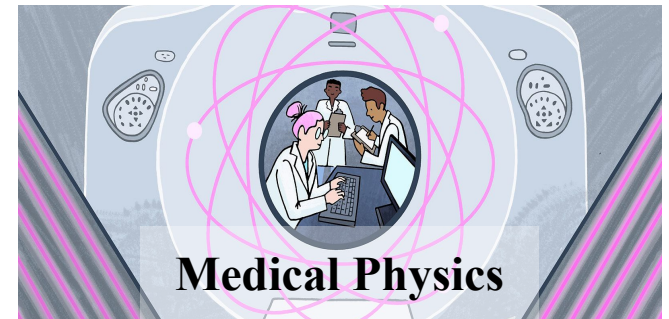
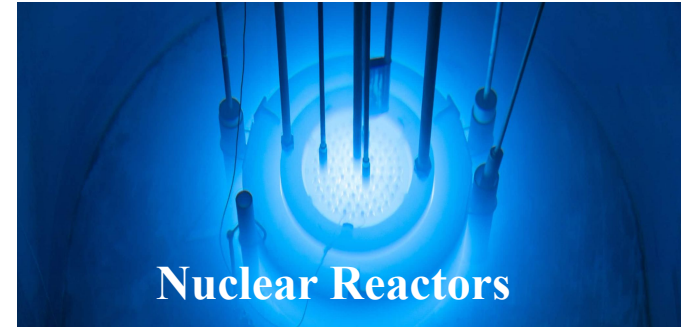
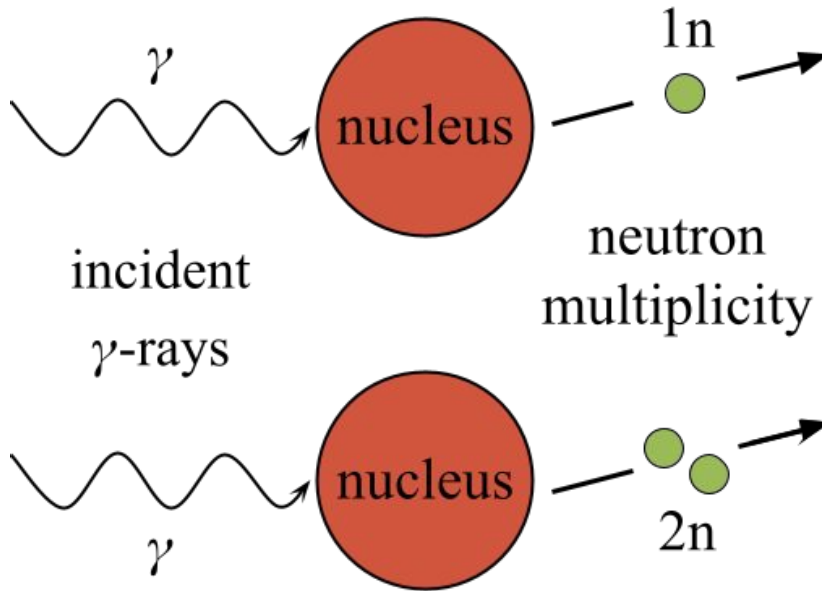
Texas A&M University

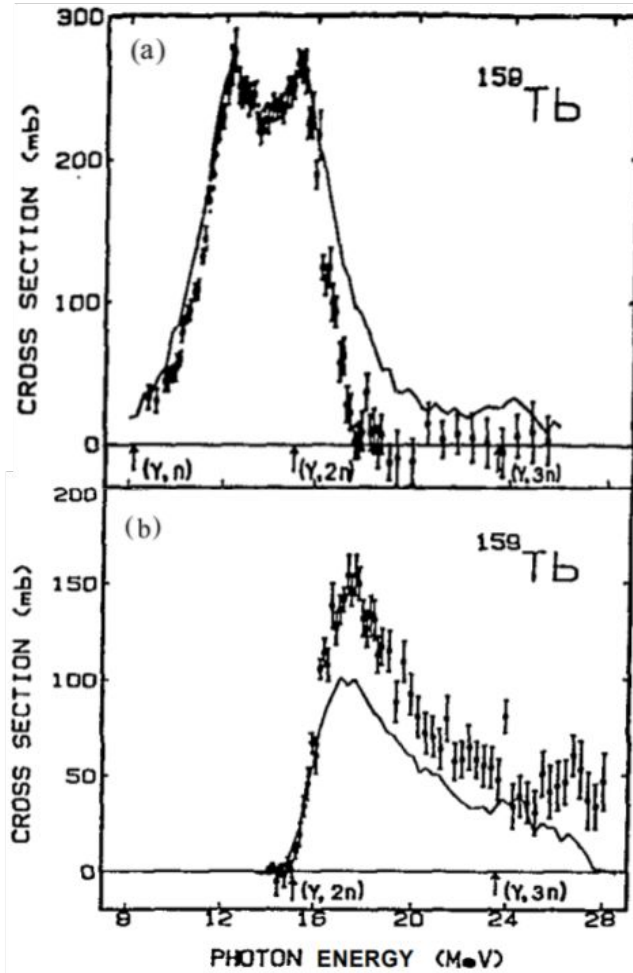
Department of Physics and Astronomy

Cyclotron Institute



Photoneutron Reaction Diagram



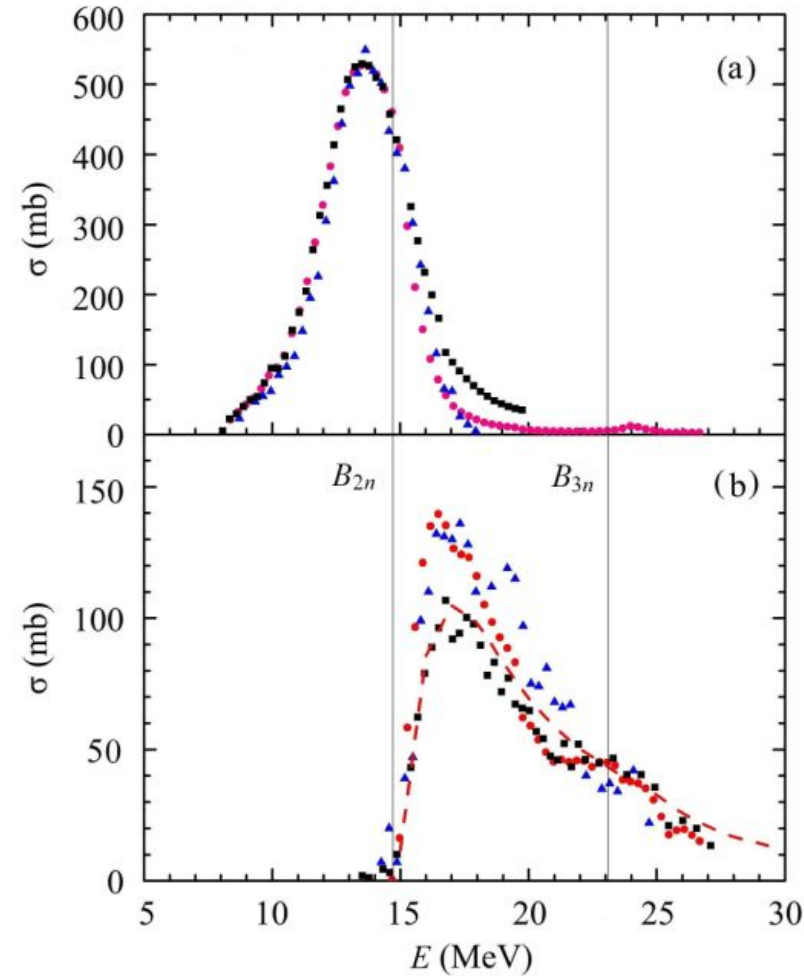


LEFT

(a) $^{159}\text{Tb}(\gamma, 1n)$ and
(b) $^{159}\text{Tb}(\gamma, 2n)$ cross
sections from Saclay
(solid line) and
Livermore (data
points).

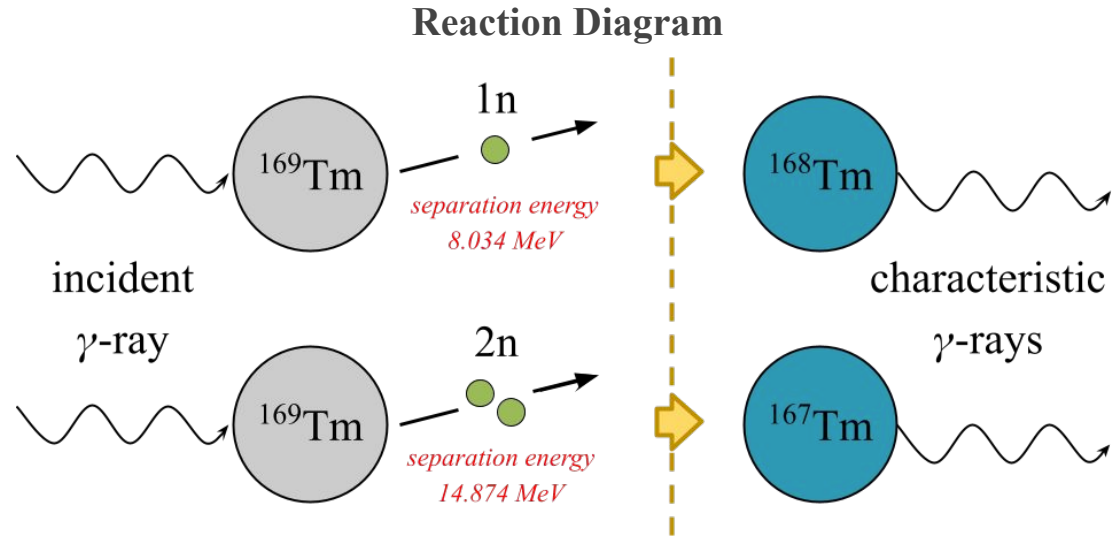
RIGHT

(a) $^{197}\text{Au}(\gamma, 1n)$ and (b)
 $^{197}\text{Au}(\gamma, 2n)$ cross
sections from Saclay
(squares) and
Livermore (triangles),
compared with
evaluated data (circles
and dotted line).



Benefits of Thulium:

- Monoisotopic (*no random byproducts to sift through*)
- Reaction products are unstable nuclei with well-understood decays of characteristic γ -rays.

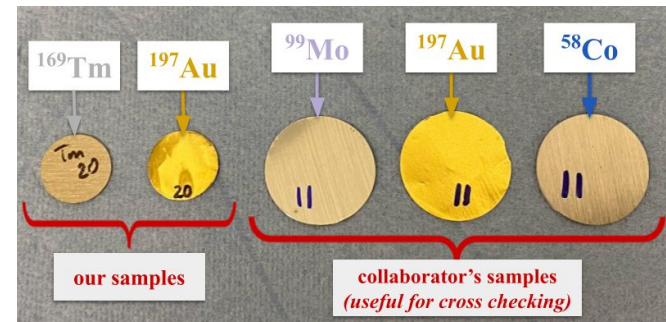
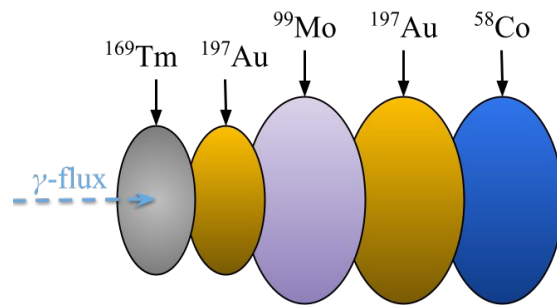
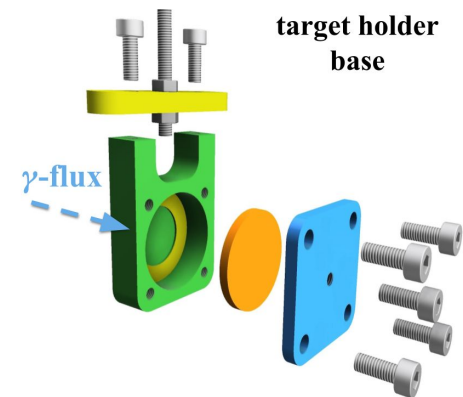
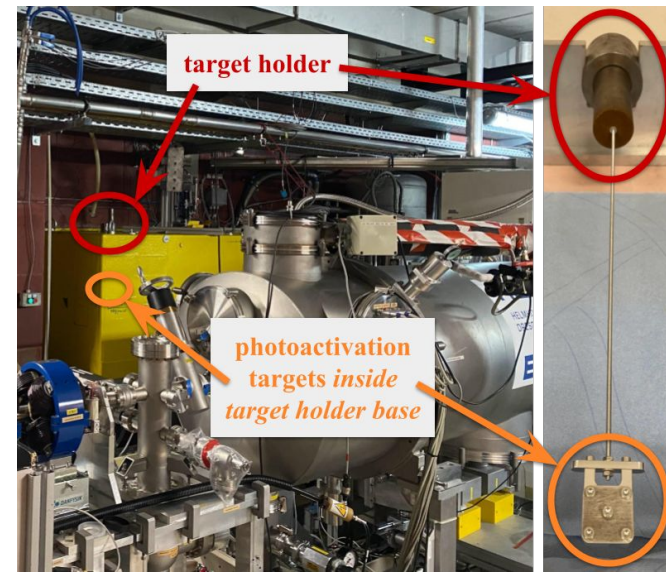
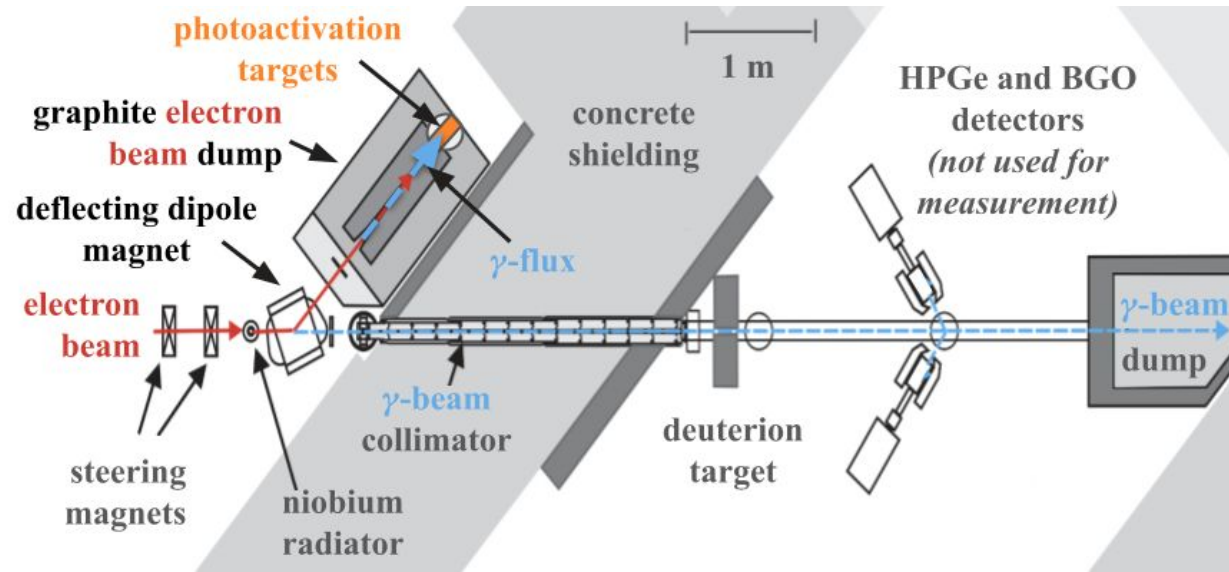


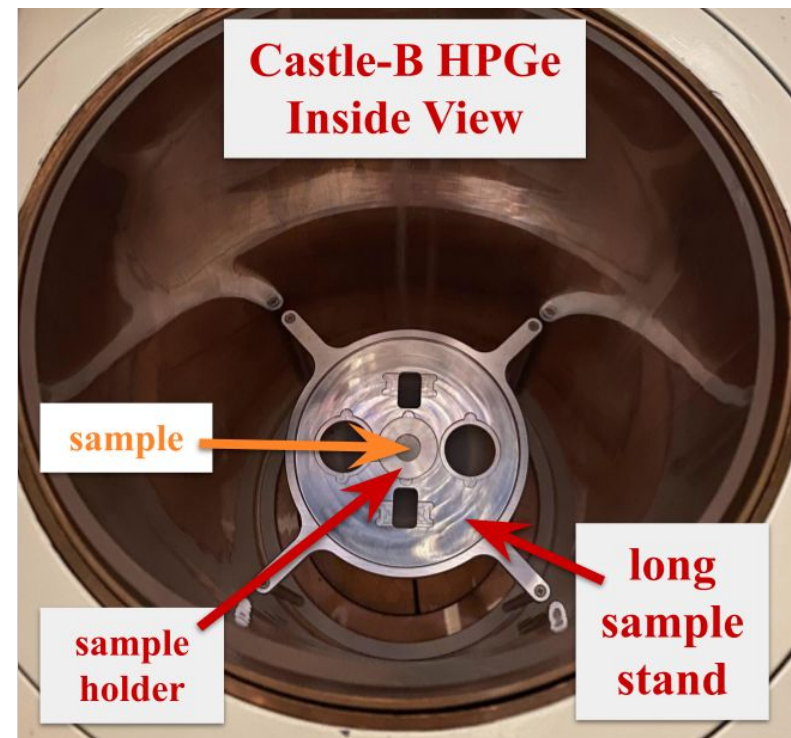
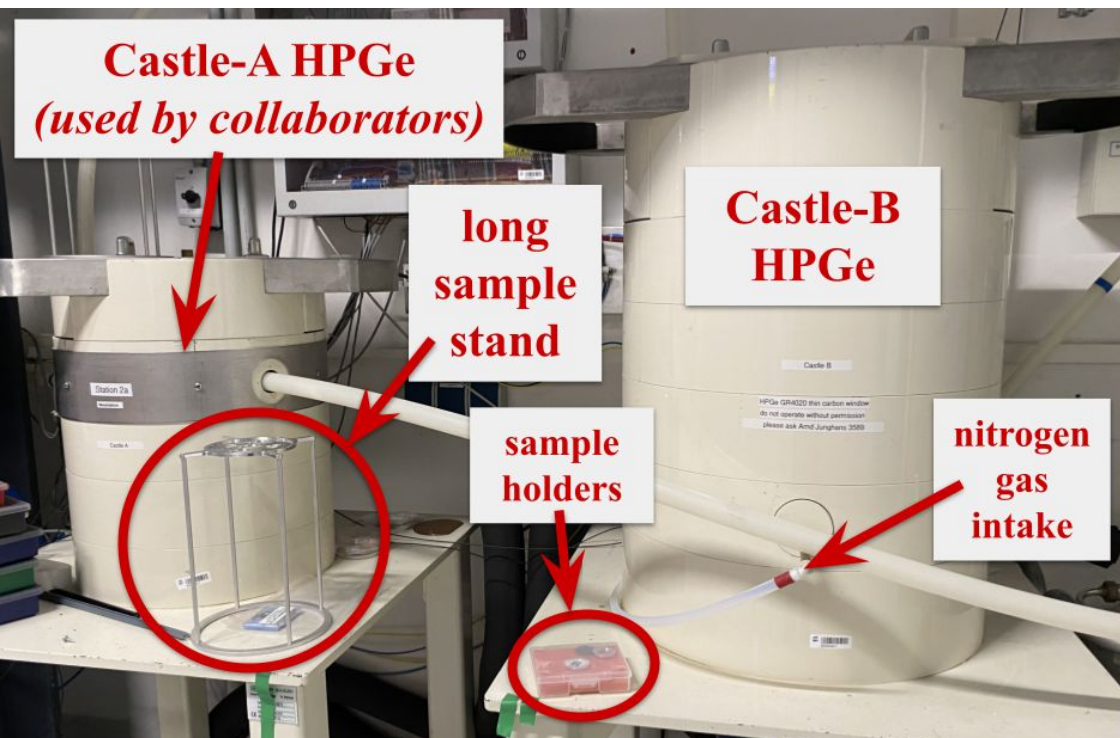
$$N_{act} = N_{tar} \int_{E_{th}}^{E_0} \sigma_{\gamma,n}(E, E_0) \Phi_{\gamma}(E, E_0) dE$$

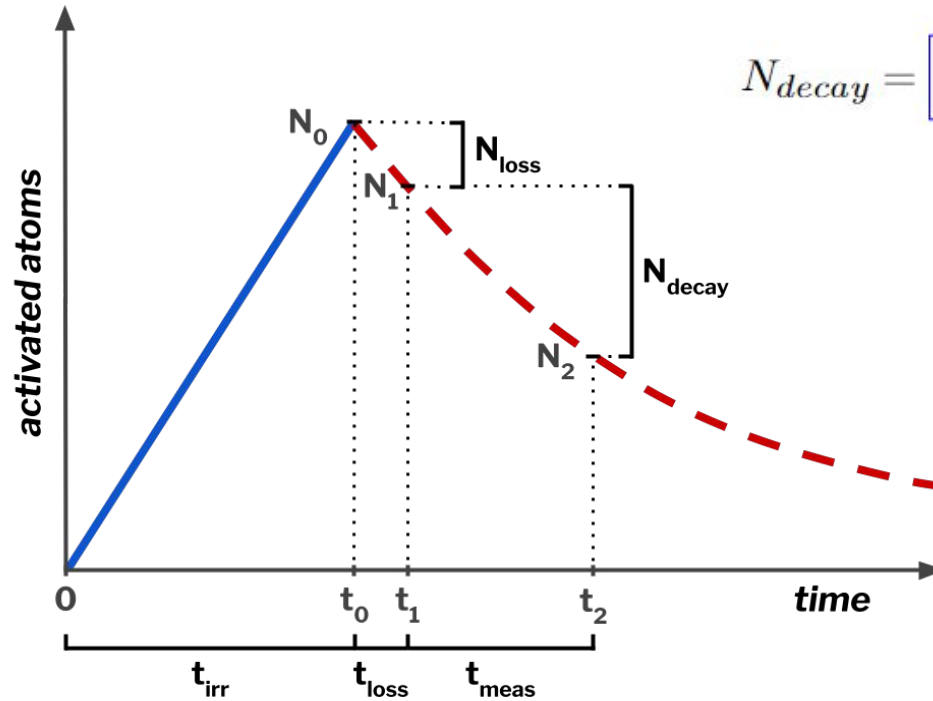
N_{act}	activated atoms	Φ_{γ}	absolute photon fluence
N_{tar}	target atoms	$\sigma_{\gamma,n}$	photoneutron cross section

Activation equation (left) connects the number activated target atoms (a measurable quantity) to the cross section.

γ ELBE Facility: Irradiation







$$N_{decay} = \underbrace{\frac{P}{\lambda}(1 - e^{-\lambda t_{irr}})}_{\text{activation}} \underbrace{e^{-\lambda t_{loss}}}_{\text{decay loss}} \underbrace{(1 - e^{-\lambda t_{meas}})}_{\text{measurement}}$$

t_{irr} irradiation time t_{meas} measurement time
 t_{loss} time from irradiation end to measurement start
 N_{decay} activated atoms which decayed during t_{meas}

$$N_{decay} = N_{meas} / \eta(E_\gamma)$$

N_{meas} total corrected counts during t_{meas}

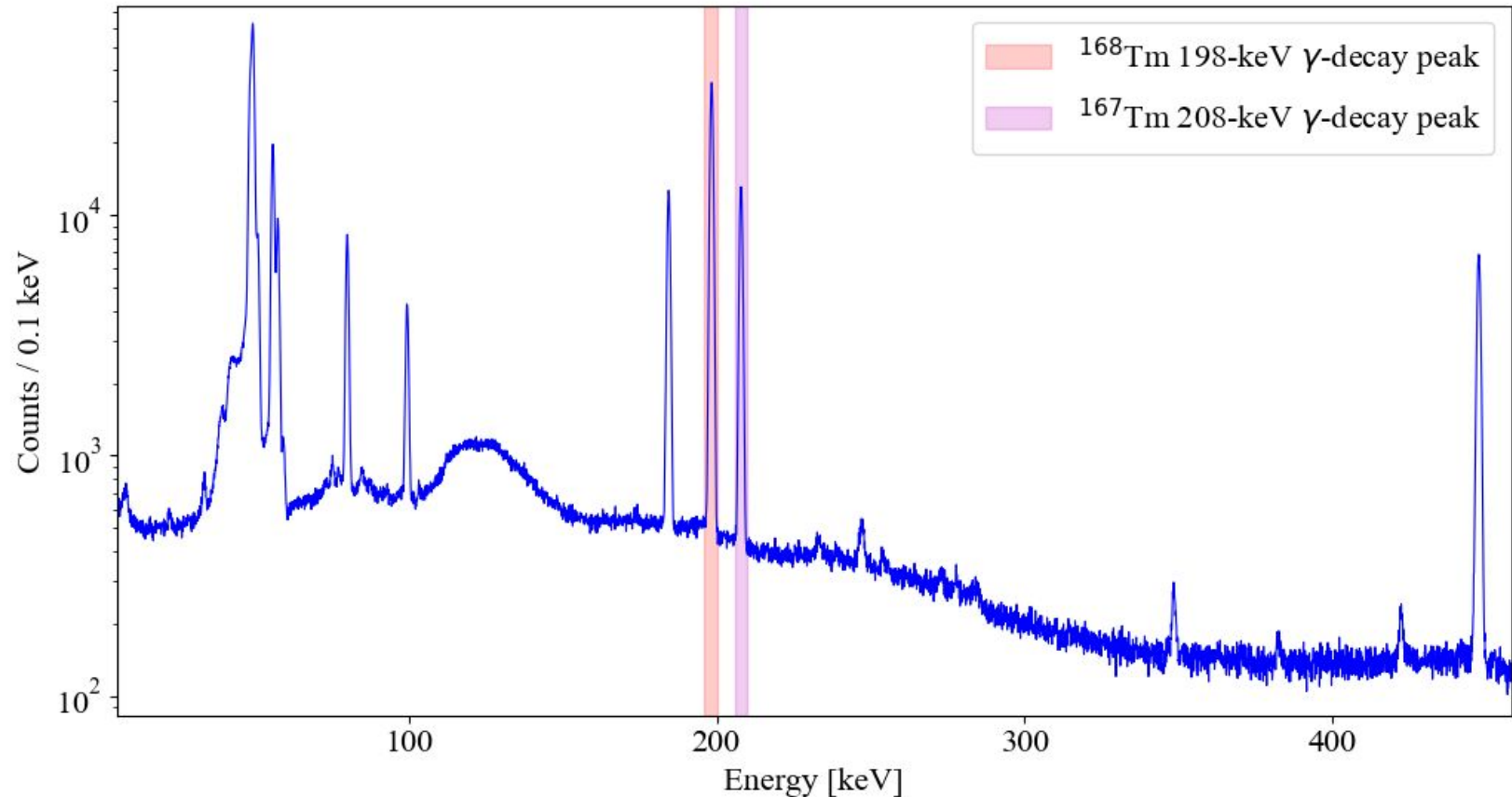
η total detector efficiency at γ -energy E_γ

End Result

$$N_{act}(t) = \frac{P}{\lambda}(1 - e^{-\lambda t}), \quad 0 \leq t \leq t_{irr}$$

P production rate λ decay constant

$$\sigma_{\gamma,n}(E_{\gamma,i-1}, E_{\gamma,i}) = \frac{N_{act}(E_{\gamma,i}) - N_{act}(E_{\gamma,i-1})}{N_{tar} \Phi_\gamma(E_{\gamma,i-1}, E_{\gamma,i})}$$



^{169}Tm target spectrum, irradiated at 18-MeV endpoint-energy γ -rays. Prominent 198-keV and 208-keV decay channel peaks from $^{169}\text{Tm}(\gamma,1n)^{168}\text{Tm}$ and $^{169}\text{Tm}(\gamma,2n)^{167}\text{Tm}$, respectively.

Thanks!
Questions?



Eq. 1 shows the differential cross section for the creation of a Bremsstrahlung with an energy E_γ in a thin target, assuming incident electron with total energy E_0 (our endpoint energy).

$$\frac{d\sigma}{dE_\gamma} = \frac{2\alpha Z^2 r_e^2}{E_\gamma} \left[\left(\frac{E_0^2 + E_e^2}{E_0^2} - \frac{2E_e}{3E_0} \right) \left(\ln M + 1 - \frac{2}{b} \arctan b \right) \right. \\ \left. + \frac{E_e}{E_0} \left(\frac{2}{b^2} \ln(1+b^2) + \frac{4}{3} \frac{2-b^2}{b^3} \arctan b - \frac{8}{3b^2} + \frac{2}{9} \right) \right] \quad (1)$$
$$M = \left[\left(\frac{m_e E_\gamma}{2E_0 E_e} \right)^2 + \frac{Z^{2/3}}{C^2} \right]^{-1} \quad b = \frac{2E_0 E_e Z^{1/3}}{C m_e E_\gamma}$$

