

BUT...

β is relative to the earth

t is not! t is the twin time

The twin is accelerating $\Rightarrow t$ is not measured in an inertial frame. It's measured in a frame where the twin is always at rest. \Rightarrow

$t = \text{Twin Wrist Watch Frame} !!!$

Called "Proper Time", and denoted τ .

Proper time is our first Lorentz Invariant,

meaning it is independent of our choice of frame.

Thus, the $\{ \}$ equation above really should have been written:

$$\left\{ \frac{d\beta}{d\tau} = \frac{g}{c} (1 - \beta^2) \right\}$$

$$\text{Then } \frac{d\beta}{dt} = \frac{d\beta}{dr} \frac{dr}{dt} = \left[\frac{g}{c} (1-\beta)^2 \right] \left(\frac{1}{\gamma} \right) = \frac{g}{c} (1-\beta^2)^{3/2}$$

↑
↑
 EARTH TIME MOVING CLOCK FACTOR

$$\Rightarrow \frac{d\beta}{(1-\beta^2)^{3/2}} = \frac{g}{c} dt \Rightarrow \int_0^\beta \frac{d\beta}{(1-\beta^2)^{3/2}} = \int_0^t \frac{g}{c} dt = \frac{g}{c} t$$

$= \frac{\beta}{\sqrt{1-\beta^2}}$ from my Int Tables

$$\Rightarrow \frac{\beta^2}{1-\beta^2} = \left(\frac{g}{c} t \right)^2$$

$$\Rightarrow \beta = \frac{\frac{g}{c} t}{\sqrt{1 + \left(\frac{g}{c} t \right)^2}}$$

with "t" measured in the (inertial) earth frame.

$$x(t) = c \int_0^t \beta(t) dt = c \frac{c}{g} \int_0^t \frac{\frac{g}{c} t}{\sqrt{1 + \left(\frac{g}{c} t \right)^2}} \frac{g}{c} dt$$

$$= \frac{c^2}{g} \int_0^{\frac{g}{c} t} \frac{y dy}{\sqrt{1+y^2}} = \frac{c^2}{g} \left[\sqrt{1+y^2} \right]_{y=0}^{\frac{g}{c} t}$$

$$= \frac{c^2}{g} \left[\sqrt{1 + \left(\frac{g}{c} t \right)^2} - 1 \right]$$

SHORT TIME LIMIT

$$= \frac{c^2}{g} \left[1 + \frac{1}{2} \frac{g^2 t^2}{c^2} - 1 \right] = \frac{1}{2} g t^2 \checkmark$$

A simplification:

"RAPIDITY" y with $\gamma = \cosh(y)$

$$\gamma^2 - \beta^2 \gamma^2 = (1 - \beta^2) \gamma^2 = \frac{1}{\gamma^2} \gamma^2 = 1 \text{ and}$$

$$\cosh^2 y - \sinh^2 y = 1 \Rightarrow$$

$$\beta \gamma = \sinh(y)$$

$$\beta = \frac{\beta \gamma}{\gamma} = \frac{\sinh(y)}{\cosh(y)} \Rightarrow$$

$$\beta = \tanh(y)$$

Velocity Addition, version 2:

$$\begin{pmatrix} x'' \\ x' \end{pmatrix} = \begin{pmatrix} \gamma_2 & -\beta_2 \gamma_2 \\ -\beta_2 \gamma_2 & \gamma_2 \end{pmatrix} \begin{pmatrix} \gamma_1 & -\beta_1 \gamma_1 \\ -\beta_1 \gamma_1 & \gamma_1 \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \end{pmatrix}$$

$$= \begin{pmatrix} \cosh y_2 & -\sinh y_2 \\ -\sinh y_2 & \cosh y_2 \end{pmatrix} \begin{pmatrix} \cosh y_1 & -\sinh y_1 \\ -\sinh y_1 & \cosh y_1 \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \end{pmatrix}$$

$$= \begin{pmatrix} \cosh y_2 \cosh y_1 + \sinh y_2 \sinh y_1 & -\dots \\ -(\sinh y_2 \cosh y_1 + \cosh y_2 \sinh y_1) & \dots \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \end{pmatrix}$$

$$= \begin{pmatrix} \cosh(y_1+y_2) & -\sinh(y_1+y_2) \\ -\sinh(y_1+y_2) & \cosh(y_1+y_2) \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \end{pmatrix} \Rightarrow$$

$$\boxed{y_3 = y_1 + y_2} \quad \text{LINEAR ADDITION!}$$

"Comfortable Twin, try 2"

$$dy = \frac{g}{c} d\tau \Rightarrow y = \frac{g}{c} \tau \Rightarrow \boxed{\beta = \tanh y = \tanh\left(\frac{g\tau}{c}\right)}$$

Same as before.
Lots easier!

$$\frac{dt}{d\tau} = \gamma = \cosh(y) = \cosh\left(\frac{g}{c}\tau\right) \Rightarrow$$

$$dt = \cosh\left(\frac{g}{c}\tau\right) d\tau \Rightarrow \left\{ t = \frac{c}{g} \sinh\left(\frac{g}{c}\tau\right) \right\}$$

$$\Rightarrow \beta(t) = \tanh\left[\frac{g}{c} \frac{c}{g} \sinh^{-1}\left(\frac{gt}{c}\right)\right]$$

$$\text{But } \tanh(\alpha) = \frac{\sinh(\alpha)}{\cosh(\alpha)} = \frac{\sinh(\alpha)}{\sqrt{1 + \sinh^2(\alpha)}} \Rightarrow$$

$$\left\{ \beta(t) = \frac{\frac{g}{c}t}{\sqrt{1 + \left(\frac{g}{c}t\right)^2}} \right\}$$

Same as before,
no Int tables needed.

Now we can find $x(t)$ by integrating:

Note: Must do $\int dt$, not $\int dx$, because $\beta = \frac{1}{c} \frac{dx}{dt}$

$$\begin{aligned}x(t) &= c \int_0^t \beta(t) dt = c \frac{c}{g} \int_0^t \frac{\frac{g}{c} t}{\sqrt{1 + \left(\frac{gt}{c}\right)^2}} \frac{g}{c} dt \\&= \frac{c^2}{g} \int_0^{\frac{gt}{c}} \frac{y}{\sqrt{1+y^2}} dy = \frac{c^2}{g} \left[\sqrt{1+y^2} \right]_{y=0}^{\frac{gt}{c}} \Rightarrow\end{aligned}$$

$$x(t) = \frac{c^2}{g} \left[\sqrt{1 + \left(\frac{gt}{c}\right)^2} - 1 \right]$$

$$\text{Short time limit} = \frac{c^2}{g} \left[\left(1 + \frac{1}{2} \frac{g^2 t^2}{c^2} \right) - 1 \right]$$

$$= \frac{1}{2} g t^2 = \text{Freshman Physics!}$$

Back to Lorentz transformations:

Λ^{μ}_{ν} = 4x4 matrix that changes frames,
with μ, ν Greek indices that span 0, 1, 2, 3

(Roman letters i, j, \dots only span 1, 2, 3 and give vectors only in space, not spacetime)

$$x^{\mu'} = \Lambda^{\mu}_{\nu} x^{\nu}$$

Einstein summation notation

BUT only sum an upper index with a lower index, or you've made a mistake!

To shift an index up or down, use the:

METRIC TENSOR

$$g_{\mu\nu} = g^{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & +1 & 0 \\ 0 & 0 & 0 & +1 \end{pmatrix}$$

This is $g^{\mu\nu}$ in Griffiths and in general relativity.

Nuclear and High Energy use: $\begin{pmatrix} +1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$

$$x_{\mu} = g_{\mu\nu} x^{\nu} \quad \text{gives}$$

LOWER
UPPER

NOTE!

$$x_0 = -ct$$

$$x_1 = x$$

$$x_2 = y$$

$$x_3 = z$$

WHY?

DEFINITIONS:

A CONTRAVARIANT 4-VECTOR A^{μ} :

Any A^{μ} that transforms like x^{μ} under Lorentz transformations

A COVARIANT 4-VECTOR A_{μ} :

Any A_{μ} that transforms like x_{μ} under Lorentz transformations.

Side note: 4-vectors are not written with arrows.

$$\vec{A} = \begin{pmatrix} A_x \\ A_y \\ A_z \end{pmatrix} \text{ is the spatial 3-vector part of } A^{\mu} \text{ or } A_{\mu}.$$

Dot product $A \cdot B = A^{\mu} B_{\mu} = A_{\nu} B^{\nu}$

$= -A^0 B^0 + A^1 B^1 + A^2 B^2 + A^3 B^3$

is Lorentz Invariant

Proof (for a boost along x):

$$\begin{aligned}A^{\mu'} B_{\mu'} &= -(\gamma A^0 - \beta \gamma A^1)(\gamma B^0 - \beta \gamma B^1) \\ &\quad + (\gamma A^1 - \beta \gamma A^0)(\gamma B^1 - \beta \gamma B^0) + A^2 B^2 + A^3 B^3 \\ &= -\gamma^2 [A^0 B^0 - \beta A^0 B^1 - \beta A^1 B^0 + \beta^2 A^1 B^1] \\ &\quad + \gamma^2 [A^1 B^1 - \beta A^1 B^0 - \beta A^0 B^1 + \beta^2 A^0 B^0] + A^2 B^2 + A^3 B^3 \\ &= -\gamma^2 (1 - \beta^2) A^0 B^0 + \gamma^2 (1 - \beta^2) A^1 B^1 + A^2 B^2 + A^3 B^3 \\ &= -A^0 B^0 + A^1 B^1 + A^2 B^2 + A^3 B^3 \quad \checkmark\end{aligned}$$

γ^2 - SQUARED!

$= x' \cdot x = x^\mu x_\mu$ Plays the role of
the vector norm squared.

Linearity of Lorentz transformations means that, if
 $x_{(1)}$ and $x_{(2)}$ are two points in spacetime,

POINT LABELS
NOT LORENTZ INDICES

then $\Delta x = x_2 - x_1$ is also a 4-vector and

$I = \Delta x^\mu \Delta x_\mu = -c^2 (\Delta t)^2 + \Delta x^2 + \Delta y^2 + \Delta z^2$ is invariant
Griffiths notation for the "Space-time Interval"