

PHYS 606 – Spring 2017 – Homework III – Solution

Problem [1]

(a) Induction: eqs (2) + (3) obviously true for $n=1$; suppose they are true for $n-1$

$$\text{then } \left. \begin{aligned} [F, G^n] &= G^{n-1}[F, G] + [F, G^{n-1}]G = G^{n-1}[F, G] + (n-2)G^{n-2}[F, G]G \\ [F^n, G] &= F^{n-1}[F, G] + [F^{n-1}, G]F = F^{n-1}[F, G] + (n-2)F^{n-2}[F, G]F \end{aligned} \right\} \begin{array}{l} \text{commutes w/ } [F, G] \\ \text{Product rule} \end{array}$$

$$= \begin{cases} (n-1)G^{n-1}[F, G] \\ (n-1)F^{n-1}[F, G] \end{cases}$$

Direct calculation is also possible.

$$\begin{aligned} (b) [F, [G, H]] &= [F, GH] - [F, HG] = G[F, H] + [F, G]H - H[F, G] - [F, H]G \\ &= [G, [F, H]] - [H, [F, G]] = -[H, [F, G]] - [G, [H, F]] \end{aligned}$$

Problem [2]

$$[\hat{F}_{\vec{r}}(\vec{r}), G_{\vec{p}}(\vec{p})] f(\vec{r}) = (2\pi\hbar)^{-3/2} \int_{\mathbb{R}^3} [F(\vec{r}), G(\vec{p})] \hat{f}(\vec{p}) e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} d^3 p$$

$$= (2\pi\hbar)^{-3/2} \int_{\mathbb{R}^3} \left(\underbrace{F(\vec{r}) \hat{f}(\vec{p}) G(\vec{p})}_{\substack{\text{power series of} \\ G \text{ has acted on} \\ \text{phase factor}}} e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} - \underbrace{G(-i\hbar \nabla_r) \hat{f}(\vec{p}) F(-i\hbar \nabla_r)}_{\substack{\text{power series} \\ \text{of } F \text{ acting} \\ \text{on phase factor}}} e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} \right) d^3 p$$

$$= (2\pi\hbar)^{-3/2} \int_{\mathbb{R}^3} \left(G(\vec{p}) \hat{f}(\vec{p}) \underbrace{F(-i\hbar \nabla_r)}_{\substack{\text{power series} \\ \text{of } F \text{ acting} \\ \text{on phase factor}}} e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} - \hat{f}(\vec{p}) \underbrace{F(-i\hbar \nabla_r)}_{\substack{\text{power series} \\ \text{of } F \text{ acting} \\ \text{on phase factor}}} G(\vec{p}) e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} \right) d^3 p$$

$$= (2\pi\hbar)^{-3/2} \int_{\mathbb{R}^3} \left(\underbrace{F(+i\hbar \nabla_r) G(\vec{p}) \hat{f}(\vec{p})}_{\substack{\text{multiple partial integrations} \\ \text{according to the power series} \\ \text{of } F; \text{ boundary terms vanish} \\ \text{for sufficiently fast falling} \\ \text{functions.}}} e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} - G(\vec{p}) \underbrace{F(+i\hbar \nabla_r) \hat{f}(\vec{p})}_{\substack{\text{partial} \\ \text{integrations}}} e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} \right) d^3 p$$

$$= (2\pi\hbar)^{-3/2} \int_{\mathbb{R}^3} [\hat{F}_{\vec{r}}(\vec{r}), G_{\vec{p}}(\vec{p})] \hat{f}(\vec{p}) e^{\frac{i}{\hbar} \vec{r} \cdot \vec{p}} d^3 p$$

$$[3] \quad \frac{d}{dt} \langle p \rangle = \langle F \rangle = 0 \quad (\text{free particle}) \Rightarrow \langle p \rangle(t) = \langle p \rangle(0) = \text{const.}$$

(Ehrenfest)

$$\frac{d}{dt} \langle x \rangle = \frac{\langle p \rangle}{m} = \frac{p_0}{m}$$

$\underbrace{\quad}_{=: p_0} \text{ short notation}$

$$\Rightarrow \langle x \rangle(t) = \frac{p_0}{m} t + x_0 \quad x_0 = \langle x \rangle(0)$$

$$\frac{d}{dt} \langle p^2 \rangle = \frac{1}{i\hbar} \langle [p^2, T] \rangle = 0 \quad \Rightarrow \langle p^2 \rangle(t) = \langle p^2 \rangle(0) = \text{const.}$$

$$\Rightarrow (\Delta p)^2(t) = \langle (p - \langle p \rangle)^2 \rangle = \langle p^2 \rangle(t) - \langle p \rangle(t)^2 = \langle p^2 \rangle(0) - p_0^2 = (\Delta p)^2(0) = \text{const.}$$

$$\begin{aligned} \frac{d}{dt} \langle x^2 \rangle &= \frac{1}{i\hbar} \langle [x^2, \frac{p^2}{2m}] \rangle = \frac{1}{i\hbar} \langle p[x^2, \frac{p}{2m}] + [x^2, \frac{p}{2m}]p \rangle \\ &= \frac{1}{i\hbar} \langle px[x, p] + p[x, p]x + x[x, p]p + [x, p]xp \rangle \frac{1}{2m} \\ &= \frac{1}{m} \langle px + xp \rangle \end{aligned}$$

$$\frac{d}{dt} \langle px \rangle = \frac{1}{i\hbar} \langle [px, T] \rangle = 2T = \text{const.} = \frac{d}{dt} \langle xp \rangle$$

$$\Rightarrow \langle px + xp \rangle(t) = \frac{2\langle p^2 \rangle(0)}{m} t + \langle px + xp \rangle(0)$$

$$\Rightarrow \langle x^2 \rangle = \frac{\langle p^2 \rangle(0)}{m^2} t^2 + \frac{1}{m} \langle px + xp \rangle(0) t + \langle x^2 \rangle(0)$$

$$\begin{aligned} \Rightarrow (\Delta x)^2(t) &= \langle x^2 \rangle(t) - \langle x \rangle(t)^2 = \frac{\langle p^2 \rangle(0) - p_0^2}{m^2} t^2 + \frac{1}{m} (\langle px + xp \rangle(0) - 2x_0 p_0) \\ &\quad + \langle x^2 \rangle(0) - x_0^2 \\ &= \frac{(\Delta p)^2(0)}{m^2} t^2 + \frac{2}{m} \left(\frac{1}{2} \langle px + xp \rangle(0) - \langle p \rangle(0) \langle x \rangle(0) \right) t + (\Delta x)^2(0) \end{aligned}$$

$$[4] (a) \text{ Gauss: } (\Delta x)^2(0) = \sigma^2, \quad (\Delta p)^2(0) = \frac{\hbar^2}{4\sigma^2} \quad (\text{HW I, [3.]})$$

$$\langle x \rangle(0) = x_0, \quad \langle p \rangle(0) = \hbar k_0 = p_0$$

$$\begin{aligned} \langle xp \rangle(0) &= \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} e^{-\frac{(x-x_0)^2}{4\sigma^2}} e^{-\frac{i}{\hbar} p_0 x} x \left(-i\hbar \frac{d}{dx}\right) e^{-\frac{(x-x_0)^2}{4\sigma^2}} e^{+\frac{i}{\hbar} p_0 x} dx \\ &= \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} e^{-\frac{(x-x_0)^2}{2\sigma^2}} x \left(p_0 + i\hbar \frac{x-x_0}{2\sigma^2}\right) dx \end{aligned}$$

$$= \frac{1}{\sqrt{2\pi}\sigma} \int_{u=x-x_0} e^{-\frac{u^2}{2\sigma^2}} \left(u p_0 + i\hbar \frac{u^2}{2\sigma^2} + x_0 p_0 + i\hbar x_0 \frac{u}{2\sigma^2} \right) du$$

$$= x_0 p_0 + i\hbar \frac{\sigma^2}{2\sigma^2} = x_0 p_0 + \frac{1}{2} i\hbar$$

$\langle p_x \rangle(0) =$ same as above + term with $(-i\hbar \frac{d}{dx})$ acting on x

$$= x_0 p_0 + \frac{1}{2} i\hbar + \frac{1}{\sqrt{2\pi}\sigma} \int_{\mathbb{R}} e^{-\frac{u^2}{2\sigma^2}} (-i\hbar) = x_0 p_0 - \frac{1}{2} i\hbar$$

$$\Rightarrow \frac{1}{2} \langle p_x + x p \rangle(0) = x_0 p_0$$

$$\Rightarrow (\Delta x)^2(t) = \frac{\hbar^2}{4\sigma^2 m^2} t^2 + \sigma^2 \quad \text{and} \quad (\Delta p)^2(t) = \frac{\hbar^2}{4\sigma^2}$$

$$(b) (\Delta p)^2(t) = \frac{1}{\sqrt{2\pi}\hat{\sigma}} \int e^{-\frac{(p-p_0)^2}{4\hat{\sigma}^2}} \left| e^{-\frac{i}{\hbar}(p-p_0)x_0} e^{-i\omega(p)t} \right|^2 (p-p_0)^2 dp \quad \left(\hat{\sigma} = \frac{\hbar}{2\sigma} \right)$$

$$= (\Delta p)^2(0) = \hat{\sigma}^2 = \frac{\hbar^2}{4\sigma^2}$$

$$\text{HW II, [1](a): } |\psi(x,t)|^2 = \frac{1}{\sqrt{2\pi}} \frac{\sigma}{\sqrt{\sigma^4 + \frac{\hbar^2 t^2}{4m^2}}} e^{-\frac{(x-x_0 - v_{gr}t)^2 \sigma^2}{2(\sigma^4 + \frac{\hbar^2 t^2}{4m^2})}} \quad v_{gr} = \frac{p_0}{m}$$

$$\langle x \rangle(t) = \frac{1}{\sqrt{2\pi}} \frac{\sigma}{\sqrt{\sigma^4 + \frac{\hbar^2 t^2}{4m^2}}} \int_{\mathbb{R}} du \underbrace{(x_0 + v_{gr}t + u)}_{=x} e^{-\frac{u^2 \sigma^2}{2(\sigma^4 + \frac{\hbar^2 t^2}{4m^2})}} = x_0 + v_{gr}t$$

$x - x_0 - v_{gr}t = u$

$$\langle x^2 \rangle(t) = \frac{1}{\sqrt{2\pi}} \frac{\sigma}{\sqrt{\sigma^4 + \frac{\hbar^2 t^2}{4m^2}}} \int_{\mathbb{R}} du \left((x_0 + v_{gr}t)^2 + 2u(x_0 + v_{gr}t) + u^2 \right) e^{-\frac{u^2 \sigma^2}{2(\sigma^4 + \frac{\hbar^2 t^2}{4m^2})}}$$

$$= (x_0 + v_{gr}t)^2 + \frac{\sigma^4 + \frac{\hbar^2 t^2}{4m^2}}{\sigma^2}$$

$$\Rightarrow \langle (x - \langle x \rangle)^2 \rangle = \langle x^2 \rangle(t) - (\langle x \rangle(t))^2 = \sigma^2 + \frac{\hbar^2}{4\sigma^2 m^2} t^2$$