

### Homework Assignment #3

(Due Date: Thursday, October 02, 05:30 pm, in class)

#### 3.1 Free and In-Medium Nucleon-Nucleon Scattering: $T$ - and $G$ -matrix (3+3+3+2 pts.)

In class the Lippmann-Schwinger equation (LSE) for the  $T$ -matrix for the scattering of 2 nucleons in free space has been introduced,

$$T(E; \vec{k}', \vec{k}) = V(\vec{k}', \vec{k}) + \int \frac{d^3p}{(2\pi)^3} V(\vec{k}', \vec{p}) \frac{1}{E - p^2/m + i\eta} T(E; \vec{p}, \vec{k}), \quad (1)$$

where  $\pm\vec{k}$  and  $\pm\vec{k}'$  denote the relative momentum of the 2 nucleons before and after the scattering, and  $E$  is the total kinetic energy, with  $E = k^2/m = k'^2/m$ ,  $m=940$  MeV the nucleon mass and  $\eta$  infinitesimal (neglect spin-isospin except for nuclear densities).

(a) Using the partial wave expansion for both potential  $V$  and  $T$ -matrix,

$$X(\vec{k}', \vec{k}) = 4\pi \sum_{l=0}^{\infty} (2l+1) P_l(u_{k'k}) X_l(k', k), \quad (2)$$

as well as azimuthal symmetry, show that the LSE can be reduced to

$$T_l(E; k', k) = V_l(k', k) + \frac{2}{\pi} \int p^2 dp V_l(k', p) \frac{1}{E - p^2/m + i\eta} T_l(E; p, k), \quad (3)$$

in each partial wave  $l$ . In eq. (2),  $u_{k'k} = \cos(\alpha_{k'k})$  with  $\alpha_{k'k} = \angle(\vec{k}', \vec{k})$ ; make use of the following identity for the Legendre polynomials

$$\int du_{pk} P_{l'}(u_{k'p}) P_l(u_{pk}) = \frac{2\delta_{ll'}}{2l+1} P_l(u_{k'k}). \quad (4)$$

(b) Concentrate on the  $S$ -wave channel ( $l=0$ ) and approximate the  $N$ - $N$  interaction by the sum of 2 schematic potentials associated with  $\sigma$ - and  $\omega$ -meson exchange,

$$V_0^\sigma(k', k) = -\frac{g_\sigma^2}{4\pi} \frac{F_\sigma(k)F_\sigma(k')}{m_\sigma^2}, \quad V_0^\omega(k', k) = +\frac{g_\omega^2}{4\pi} \frac{F_\omega(k)F_\omega(k')}{m_\omega^2}; \quad (5)$$

hadronic formfactors,  $F(k) = \Lambda^2/(\Lambda^2 + k^2)$ , simulate the finite size of the hadrons and ensure the convergence of the 1-D LSE. Write down explicitly the first 3 terms of the Born series for the  $T$ -matrix; exploit the separability of the above potentials,  $V(k', k) = v(k')v(k)$ , to resum the  $T$ -matrix using a geometric series to obtain

$$T_0(E; k', k) = \frac{V_0(k', k)}{1 - \Pi(E)}, \quad \Pi(E) = \frac{2}{\pi} \int p^2 dp \frac{V_0(p)}{E - p^2/m + i\eta} \quad (6)$$

(c) For  $S$ -wave scattering, the total  $N$ - $N$  cross section takes the form

$$\sigma_{\text{tot}}^{l=0}(E) = 4\pi m^2 |T_0(E)|^2 . \quad (7)$$

To evaluate the “loop” function,  $\Pi(E)$ , make us of the following decomposition into real and imaginary parts:

$$\int dp \frac{f(p)}{p_0^2 - p^2 + i\eta} = PP \int_0^\infty dp \frac{f(p)}{p_0^2 - p^2} + \int dp f(p) (-i\pi) \delta(p_0^2 - p^2) \quad (8)$$

(here:  $p_0^2 = mE$ ). The principle-value ( $PP$ ) integral for the real part of  $\Pi(E)$  requires a numerical integration. To avoid numerical instabilities when integrating over the pole, use the following “regularization” trick

$$PP \int_0^\infty dp \frac{f(p)}{p_0^2 - p^2} = PP \int_0^\infty dp \frac{f(p) - f(p_0)}{p_0^2 - p^2} , \quad \text{since} \quad PP \int_0^\infty \frac{dp}{p_0^2 - p^2} = 0 \quad (9)$$

Plot the cross section in  $[\text{mb}] = [0.1 \text{ fm}^2]$  from threshold to  $E=150$  MeV using the following values for the potential parameters:  $m_\sigma=550$  MeV,  $g_\sigma=10.0$ ,  $\Lambda_\sigma=1000$  MeV,  $m_\omega=782.6$  MeV,  $g_\omega=11.65$ ,  $\Lambda_\omega=1500$  MeV; use  $\hbar c=197.33$  MeVfm.

(d) Compute the in-medium  $N$ - $N$   $G$ -matrix by implementing a Pauli-Blocking factor,  $[1 - f(\epsilon_p; \mu_N, T)]^2$ , into the integral of the LSE (3), where

$$f(\epsilon_p; \mu_N, T) = \frac{1}{\exp[(\epsilon_p - \mu_N)/T] + 1} \quad (10)$$

with  $\epsilon_p = p^2/2m$  and chemical potential  $\mu_N = k_F^2/2m$  ( $k_F$ : Fermi momentum). Plot the in-medium  $N$ - $N$   $S$ -wave cross section for  $k_F=220$  MeV with temperature  $T=1$  MeV and  $T=5$  MeV.