

**Physics 606 -- Problem Set 3**  
Due Thursday, February 19, 2009

Do the following problems from Merzbacher:

- Chapter 4: Exercises 4.11, 4.13, and 4.17
- Chapter 4: End-of-chapter Problem 3

## Chapter 4, Exercise 11

If  $\psi = c_1\psi_1 + c_2\psi_2$ , then  $\langle U\psi | U\psi \rangle = \langle \psi | \psi \rangle$ ,  
 $\langle U\psi_1 | U\psi_1 \rangle = \langle \psi_1 | \psi_1 \rangle$ ,  $\langle U\psi_2 | U\psi_2 \rangle = \langle \psi_2 | \psi_2 \rangle$

$$\langle U\psi | U\psi \rangle = c_1^* c_1 \langle U\psi_1 | U\psi_1 \rangle + c_1^* c_2 \langle U\psi_1 | U\psi_2 \rangle$$

$$+ c_2^* c_1 \langle U\psi_2 | U\psi_1 \rangle + c_2^* c_2 \langle U\psi_2 | U\psi_2 \rangle$$

$$= c_1^* c_1 \langle \psi_1 | \psi_1 \rangle + c_1^* c_2 \langle \psi_1 | \psi_2 \rangle$$

$$+ c_2^* c_1 \langle \psi_2 | \psi_1 \rangle + c_2^* c_2 \langle \psi_2 | \psi_2 \rangle$$

$$\Rightarrow \cancel{c_1^* c_2 \langle U\psi_1 | U\psi_2 \rangle + c_2^* c_1 \langle U\psi_2 | U\psi_1 \rangle} = c_1^* c_2 \langle \psi_1 | \psi_2 \rangle + c_2^* c_1 \langle \psi_2 | \psi_1 \rangle$$

$$c_1^* c_2 \langle U\psi_1 | U\psi_2 \rangle + c_2^* c_1 \langle U\psi_2 | U\psi_1 \rangle = c_1^* c_2 \langle \psi_1 | \psi_2 \rangle + c_2^* c_1 \langle \psi_2 | \psi_1 \rangle$$

$$\Rightarrow c_1^* c_2 \langle U\psi_1 | U\psi_2 \rangle + c_2^* c_1 \langle U\psi_2 | U\psi_1 \rangle^* = c_1^* c_2 \langle \psi_1 | \psi_2 \rangle + c_2^* c_1 \langle \psi_2 | \psi_1 \rangle^*$$

$$\text{Let } c_1 = c_2 = 1 \Rightarrow \text{Re} [\langle U\psi_1 | U\psi_2 \rangle] = \text{Re} [\langle \psi_1 | \psi_2 \rangle]$$

$$\text{Let } c_1 = 1, c_2 = -i \Rightarrow \text{Im} [\langle U\psi_2 | U\psi_1 \rangle] = \text{Im} [\langle \psi_2 | \psi_1 \rangle]$$

$$\Rightarrow \langle U\psi_1 | U\psi_2 \rangle = \langle \psi_1 | \psi_2 \rangle, \text{ for all } \psi_1, \psi_2$$

## Chapter 4, Exercise 13

$$D_{\xi} g(x) = g(x - \xi), \text{ by definition.}$$

Also, if  $g(x)$  is an eigenfunction of  $D_{\xi}$  with eigenvalue  $\lambda = e^{-ik\xi}$ ,

$$\text{then } D_{\xi} g(x) = e^{-ik\xi} g(x) = g(x - \xi).$$

$$\begin{aligned} \text{Let } g(x) &= e^{ikx} u(x). \text{ Then } e^{-ik\xi} g(x) = e^{ik(x-\xi)} u(x) = g(x - \xi) \\ &= e^{ik(x-\xi)} u(x - \xi) \end{aligned}$$

$\Rightarrow$  we must have  $u(x) = u(x - \xi)$ , and  $u$  must be periodic with period  $\xi$ .

# Chapter 4, Exercise 17

$$i\hbar \frac{\partial \psi}{\partial t} = H\psi \Rightarrow i\hbar \psi^* \frac{\partial \psi}{\partial t} = \psi^* H\psi$$

$$-i\hbar \frac{\partial \psi^*}{\partial t} = (H\psi)^* \Rightarrow -i\hbar \left( \frac{\partial \psi^*}{\partial t} \right) \psi = (H\psi)^* \psi$$

$$\Rightarrow i\hbar \frac{\partial}{\partial t} (\psi^* \psi) = \psi^* H\psi - (H\psi)^* \psi$$

$$\Rightarrow \frac{\partial}{\partial t} (\psi^* \psi) = \frac{1}{i\hbar} \left[ \psi^* H\psi - (H\psi)^* \psi \right]$$

Note:  $H\psi = \frac{1}{2m} \left( p - \frac{q}{c} A \right)^2 \psi + q\phi \psi = \frac{1}{2m} \left[ \frac{\hbar}{i} \vec{\nabla} - \frac{q}{c} A \right]^2 \psi + q\phi \psi$

$$= \left( -\frac{\hbar^2}{2m} \nabla^2 - \frac{\hbar q}{2mci} (\vec{\nabla} \cdot \vec{A}) - \frac{2\hbar q}{2mci} \vec{A} \cdot \vec{\nabla} + \frac{q^2}{2mc^2} A^2 \right) \psi + q\phi \psi$$

$$= -\frac{\hbar^2}{2m} \left[ \nabla^2 + \frac{q}{\hbar ci} (\vec{\nabla} \cdot \vec{A}) + \frac{2q}{\hbar ci} \vec{A} \cdot \vec{\nabla} - \frac{q^2}{\hbar^2 c^2} A^2 \right] \psi + q\phi \psi$$

Thus:  $\frac{\partial}{\partial t} (\psi^* \psi) = \frac{-\hbar}{2mci} \left\{ \psi^* \nabla^2 \psi + \psi^* \frac{q}{\hbar ci} (\vec{\nabla} \cdot \vec{A}) \psi + \psi^* \frac{2q}{\hbar ci} \vec{A} \cdot \vec{\nabla} \psi - \psi^* \frac{q^2}{\hbar^2 c^2} A^2 \psi - (\nabla^2 \psi^*) \psi + \frac{q}{\hbar ci} (\vec{\nabla} \cdot \vec{A}) \psi^* \psi + \frac{2q}{\hbar ci} (\vec{A} \cdot \vec{\nabla}) \psi^* \psi + \frac{q^2}{\hbar^2 c^2} A^2 \psi^* \psi \right\}$

Consider  $\vec{\nabla} \cdot \vec{y} = \frac{\hbar}{2mci} \left\{ \psi^* \left( \vec{\nabla} - \frac{iq}{\hbar c} \vec{A} \right) \psi - \left[ \left( \vec{\nabla} - \frac{iq}{\hbar c} \vec{A} \right) \psi \right]^* \psi \right\}$

Then  $\vec{\nabla} \cdot \vec{y} = \frac{\hbar}{2mci} \left\{ (\vec{\nabla} \psi^*) \cdot \vec{\nabla} \psi - \frac{iq}{\hbar c} (\vec{\nabla} \psi^*) \psi + \psi^* \nabla^2 \psi - \psi^* \frac{iq}{\hbar c} (\vec{\nabla} \cdot \vec{A}) \psi \right.$

$$\left. - \psi^* \frac{iq}{\hbar c} \vec{A} \cdot \vec{\nabla} \psi - (\nabla^2 \psi^*) \psi - \frac{iq}{\hbar c} (\vec{\nabla} \cdot \vec{A}) \psi^* \psi - \frac{iq}{\hbar c} (\vec{A} \cdot \vec{\nabla}) \psi^* \psi \right.$$

$$\left. - (\vec{\nabla} \psi^*) \cdot (\vec{\nabla} \psi) - \frac{iq}{\hbar c} \psi^* \vec{A} \cdot \vec{\nabla} \psi \right\}$$

$$\Rightarrow \frac{\partial}{\partial t} (\psi^* \psi) + \vec{\nabla} \cdot \vec{y} = 0 \quad \text{w/} \quad \vec{y} = \frac{\hbar}{2mci} \left\{ \psi^* \left( \vec{\nabla} - \frac{iq}{\hbar c} \vec{A} \right) \psi - \left[ \left( \vec{\nabla} - \frac{iq}{\hbar c} \vec{A} \right) \psi \right]^* \psi \right\}$$

# Chapter 4, Problem 3

a)  $\vec{A} = \frac{\Phi}{2\pi} \frac{\hat{k} \times \vec{r}}{(\hat{k} \times \vec{r})^2} \Rightarrow \vec{A} = \frac{\Phi}{2\pi} \begin{pmatrix} -\frac{y}{x^2+y^2} \\ \frac{x}{x^2+y^2} \\ 0 \end{pmatrix}$

Then  $B_x = \partial_y A_z - \partial_z A_y = 0 - 0 = 0$ .

Likewise,  $B_y = 0$ .

$$B_z = \partial_x A_y - \partial_y A_x = \frac{\Phi}{2\pi} \left[ \frac{(x^2+y^2) - x(2x)}{(x^2+y^2)^2} - (-) \frac{(x^2+y^2) - y(2y)}{(x^2+y^2)^2} \right]$$

$$= \frac{\Phi}{2\pi} \left[ \frac{(x^2+y^2) + (x^2+y^2) - 2x^2 - 2y^2}{(x^2+y^2)^2} \right] = 0 \text{ if } (x,y) \neq (0,0)$$

But if you integrate over a circle of radius  $R$  centered on the  $z$ -axis, you get:

~~$\int \vec{B} \cdot d\vec{A}$~~   $\int \vec{B} \cdot d\vec{A} = \int \vec{k} \times \vec{A} \cdot d\vec{A} = \oint \vec{A} \cdot d\vec{l} = \frac{\Phi}{2\pi R} \int dl = \Phi$ , so there must be a magnetic flux  $\Phi$  along the  $z$ -axis.

b)  $\vec{A} = A_\phi \hat{\phi}$ , where  $A_\phi = \frac{\Phi}{2\pi\rho}$ . Then we have

$$\left( \frac{\vec{p} - \frac{q}{c} \vec{A}}{2m} \right)^2 \psi_E = E \psi_E \Rightarrow \frac{1}{2m} \left( \frac{\hbar}{i\rho} \frac{\partial}{\partial \phi} - \frac{q\Phi}{2\pi\rho c} \right)^2 \psi_E = E \psi_E$$

We still need  $\psi_E$  periodic  $\Rightarrow \psi_E = \frac{1}{\sqrt{2\pi}} e^{in\phi} \Rightarrow$

$$\frac{1}{2m} \left( \frac{\hbar}{i\rho} \frac{\partial}{\partial \phi} - \frac{q\Phi}{2\pi\rho c} \right)^2 \psi_E = \left( \frac{\hbar n}{\rho} - \frac{q\Phi}{2\pi\rho c} \right)^2 \psi_E \Rightarrow$$

$$E_n = \frac{1}{2m} \left( \frac{\hbar n}{\rho} - \frac{q\Phi}{2\pi\rho c} \right)^2 = \frac{1}{2m\rho^2} \left[ \hbar^2 \frac{n^2}{\rho^2} - 2 \frac{\hbar n q \Phi}{2\pi\rho c} + \frac{q^2 \Phi^2}{4\pi^2 \rho^2 c^2} \right]$$

$$= \frac{\hbar^2 n^2}{2m\rho^2} - \frac{\hbar n q \Phi}{2m\rho c} + \frac{q^2 \Phi^2}{8\pi^2 m \rho^2 c^2}$$

Same if  $n\hbar = \frac{q\Phi}{2\pi c}$

Same final spectrum when  $\frac{q\Phi}{2\pi c} = n\hbar$ . Then  $E_n \rightarrow E_{n-n'}$