

J/ψ Production from Relativistic Heavy Ion Collisions

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Suppression of J/ψ production in ultrarelativistic heavy ion collisions is one of the most studied signals for the quark-gluon plasma formed in these collisions [1]. Since the quark-gluon plasma has a finite size, exists for a finite time, and may not be in equilibrium, it is important to study how suppression of J/ψ is affected. We have recently carried out such a study [2] using the AMPT model [3], which includes both initial partonic and final hadronic interactions.

To include the effect of plasma screening, we have determined in the AMPT model a time-dependent critical radius within which the plasma density is higher than the critical density of about 5 fm^{-3} , estimated from the critical Debye screening mass. A pair of charm-anticharm quark then cannot form a J/ψ if their separation is smaller than the critical radius after a formation time of $0.5 \text{ fm}/c$. In the quark-gluon plasma, a J/ψ may also be destroyed by collisions with gluons with a cross section of about 3 mb [4]. Furthermore, the J/ψ can be absorbed in the hadronic matter following the phase transition of the partonic matter, and the cross section is about 6 mb for absorption by baryons and 3 mb for absorption by mesons [5]. In both partonic and hadronic matters, J/ψ can also be produced from the inverse reactions of charm-anticharm quark and charm-anticharm meson annihilation.

We first obtain from the PYTHIA model [6] a midrapidity density of 1.5 charm-anticharm quark pair and $0.019 J/\psi$ in central Au+Au collisions at 200 AGeV . In the partonic phase,

the time evolution of the numbers of produced and destroyed J/ψ per unit rapidity as well as the net number of J/ψ per unit rapidity are shown in Fig. 1. The J/ψ number reaches the maximum at about $2 \text{ fm}/c$.

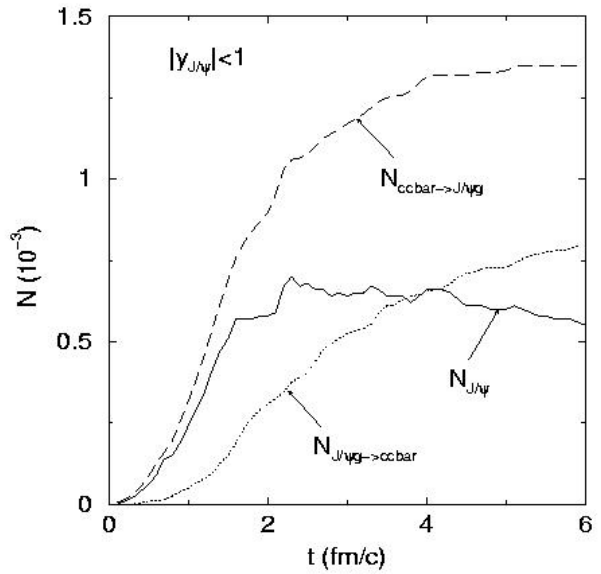


Figure 1: Time evolution of the produced number (dashed line), destroyed number (dotted line), and net number (solid line) per unit rapidity for J/ψ with $|y_{J/\psi}| < 1$ in the parton phase with color screening effect.

In the hadronic stage, the effect of charm mesons are studied by using the three charm meson masses of 1.70 , 1.87 , and 2.01 GeV to estimate the effects of D and D^* meson as well as the possible change of their masses in hadronic matter. In Fig. 2, we show the time evolution of the numbers of produced and destroyed J/ψ together with that of the net J/ψ number. For all three D meson masses, the number of produced J/ψ is always greater

than the number of destroyed J/ψ , leading to a net production of J/ψ from the hadron phase of relativistic heavy ion collisions. The time evolution of the net J/ψ number in the three cases is similar, as it is largely determined by the heavy ion collision dynamics. In all three cases, the J/ψ number saturates at about 15-20 fm/c,

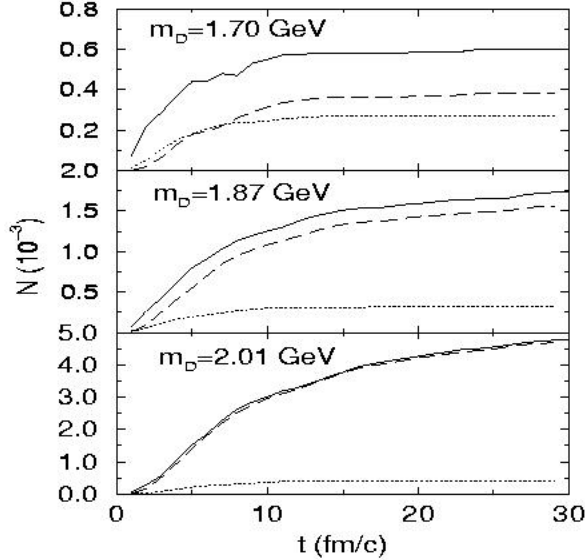


Figure 2: Time evolution of the produced number (dashed line), destroyed number (dotted line), and net number (solid line) per unit rapidity for J/ψ with $|y_{J/\psi}| < 1$ for three values of D meson mass. The color screening effect is included in the initial parton phase.

which is about the lifetime of the hadronic matter. Quantitatively, the final net number of J/ψ per unit rapidity averaged over $|y_{J/\psi}| < 1$ is about 0.0006 per event for $m_D = 1.70\text{GeV}$ and increases to 0.0017 for $m_D = 1.87\text{GeV}$ and to 0.0048 for $m_D = 2.01\text{GeV}$. In all cases, the final J/ψ number is smaller than the primary J/ψ number of 0.019. Our study thus predicts a suppression of J/ψ production in heavy ion collisions at RHIC, and this is opposite to the enhancement predicted by the kinetic formation model [7].

To predict quantitatively J/ψ production in heavy ion collisions at RHIC, we need to include the effects due to ψ' and χ_c which are also produced in the initial collisions and can decay to J/ψ . Since these particles are less bound than J/ψ , they are more likely to be dissociated and absorbed in both the quark-gluon plasma and the hadronic matter. As a result, inclusion of ψ' and χ_c is expected to lead to an even larger suppression of J/ψ production than shown here.

References

- [1] T. Matsui and H. Satz, Phys. Lett. **B178**, 416 (1986).
- [2] B. Zhang, C. M. Ko, B. A. Li, Z. W. Lin, and B. H. Sa, Phys. Rev. C **62**, 054905 (2000); *ibid.*, in press.
- [3] B. Zhang, C. M. Ko, B. A. Li, and Z. W. Lin, Phys. Rev. C **61**, 067901 (2000); Z. W. Lin, S. Pal, C. M. Ko, B. A. Li, and B. Zhang, *ibid.* **64**, 011902 (2001); Nucl. Phys. **A698**, 375c (2002).
- [4] E. V. Shuryak, Yad. fiz. **28**, 796 (1978).
- [5] S. Gavin, Nucl. Phys. A **566**, 287c (1994); W. Cassing and C. M. Ko, Phys. Lett. **B396**, 39 (1997); R. Vogt, Phys. Rep. **310**, 197 (1999).
- [6] T. Sjöstrand, Comput. Phys. Commun. **82**, 74 (1994).
- [7] R. Thews, M. Schroedter, and J. Rafelski, Phys. Rev. C **63**, 054905 (2001).