

## The $\beta^+$ Decay of $^{34}\text{Ar}$

V. E. Mayes, J. C. Hardy, V. E. Jacob, M. Sanchez-Vega, A. Azhari, C. A. Gagliardi, J. Giovinazzo<sup>a</sup>, R. G. Neilson, L. Trache, R. E. Tribble and Y. M. Xiao

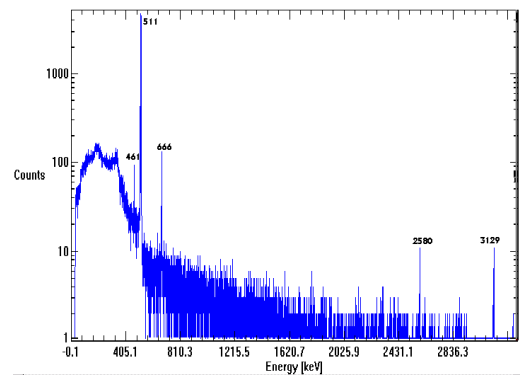
<sup>a</sup> Centre d'Etudes Nucleaires de Bordeaux- Gradignan, IN2P3/CNRS, France

We chose  $^{22}\text{Mg}$  as the first new superallowed  $\beta$  emitter to study because its calculated charge correction,  $\delta_C$ , has the very low value 0.26%, and because  $^{22}\text{Mg}$  has no  $\beta$ -decay branch feeding its daughter's ground-state; thus, its  $\beta$ -decay branching ratios can be determined entirely from the relative intensities of its  $\beta$ -delayed  $\gamma$ -rays. We have now chosen  $^{34}\text{Ar}$  ( $t_{1/2} = 0.84$  s) as the second case to study for two reasons: 1) Its calculated charge correction is 0.68%, higher than any other known case, and consequently it offers an important test [1] of the  $\delta_C$  calculations; and 2) it is the case next most tractable to experimental precision after  $^{22}\text{Mg}$ . The superallowed  $\beta$ -decay branch from  $^{34}\text{Ar}$ , which directly populates the  $^{34}\text{Cl}$  ground state, accounts for 94% of its total decay strength. This means that, unlike  $^{22}\text{Mg}$ , the branching ratios in the decay of  $^{34}\text{Ar}$  require a measurement of the *absolute* intensities of the  $\beta$ -delayed  $\gamma$ -rays from the non-superallowed transitions – a normally challenging requirement. However, because these transitions comprise less than 6% of the total decay strength from  $^{34}\text{Ar}$ , they need only be determined to a precision of 2% in order to yield 0.1% precision on the superallowed branch.

We plan to measure both the branching ratio and half-life of  $^{34}\text{Ar}$  at the cyclotron. A test run with MARS in March yielded a beam of  $^{34}\text{Ar}$  with  $\sim 95\%$  purity. We produced this beam using the reaction  $p(^{35}\text{Cl}, 2n)^{34}\text{Ar}$ , with a 30.4 MeV beam from the cyclotron bombarding a  $\text{LN}_2$ -cooled hydrogen gas target. After MARS

had been optimized, we observed  $\sim 5000$   $^{34}\text{Ar}$  ions per second at its focal plane. This beam was extracted from MARS through a Kapton window, then passed through a thin scintillator and a stack of aluminum degraders, and finally was implanted in the aluminized Mylar tape of the fast tape-transport system. (For a more complete description of the experimental arrangement, see [2].) Each collected sample was conveyed to a shielded counting station, where it was positioned between a thin plastic scintillator and a 70% HPGe detector. There, we recorded  $\beta$ - $\gamma$  coincidences event-by-event. A counting cycle consisted of a 2-s collection period, a 180-ms transfer time and a 2-s counting period. The accelerator beam was off during the transfer and counting periods.

This cycle was clock-controlled and was repeated continuously. The coincident  $\gamma$ -ray spectrum obtained in this way is shown in Figure 1. Although the counting statistics are



**Figure 1:** Coincidence ( $\gamma$ -ray spectrum at 15 cm from March run.

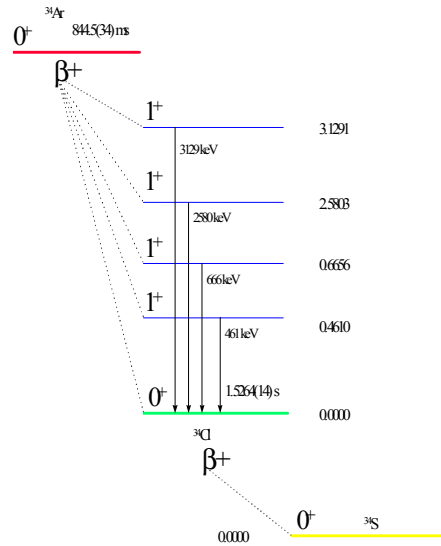
very low, there is no evidence of any radioactivity other than  $^{34}\text{Ar}$ .

The energy-level diagram for the decay of  $^{34}\text{Ar}$  is shown in Figure 2, and the current values for the known  $\gamma$ -ray intensities and  $\beta$  branching ratios are shown in the table. It is evident from the spectrum that we observe all of the  $\gamma$ -rays known to be emitted in the decay of  $^{34}\text{Ar}$ . Furthermore, the spectrum is substantially cleaner than the best spectrum obtained previously [3]. A preliminary analysis of the  $\gamma$ -ray intensities in this spectrum gives a result for the intensity of the 461 keV  $\gamma$ -ray relative to the one with 666 keV of  $0.35 \pm 0.04$ , in good agreement with the prior results[3].

In order to determine intensities for all of the  $\gamma$ -rays, it will be necessary to extend our current efficiency curve for the HPGe detector from 1800 keV to 3200 keV. This will be accomplished using a  $^{56}\text{Co}$  source. In addition, Monte Carlo simulations are currently being performed in order to determine the limits of our precision in measuring the half-life of  $^{34}\text{Ar}$  with the  $4\pi$   $\beta$ -detector [4]. Measurements of this half-life are complicated by the fact that the daughter of  $^{34}\text{Ar}$  is  $^{34}\text{Cl}$ , which has a similar short half-life ( $t_{1/2} = 1.53$  s). Any measurement

**Table 1:** Branching ratios and relative  $\gamma$ -ray intensities for  $^{34}\text{Ar}$

Excitation Energy (keV)	Relative intensity of $\gamma$ -rays (%)	Branching ratios (%)
0	0.00	$94.4 \pm 0.25$
461	$36.5 \pm 3.6$	$0.91 \pm 0.10$
666	100	$2.49 \pm 0.12$
2580	$34.5 \pm 1.0$	$0.86 \pm 0.12$
3129	$52.1 \pm 1.2$	$1.30 \pm 0.10$



**Figure 2:** Energy level diagram for the decay of  $^{34}\text{Ar}$

of  $\beta$ -particles alone will include the growth and decay of  $^{34}\text{Cl}$  along with the decay of  $^{34}\text{Ar}$ . If we cannot obtain sufficient precision with the  $4\pi$  detector, it will then be necessary to consider another method for measuring the half-life, such as the  $\beta$ - $\gamma$  coincidence method, in order to differentiate between the decay of  $^{34}\text{Ar}$  from that of  $^{34}\text{Cl}$ .

## References

- [1] J. C. Hardy *et al.*, *Progress in Research*, Cyclotron Institute, Texas A&M University (2000-2001), p. I-24.
- [2] M. Sanchez-Vega *et al.*, *Progress in Research*, Cyclotron Institute, Texas A&M University (2000-2001), p. I-27.
- [3] J. C. Hardy *et al.*, *Nucl.Phys.* **A223**, 157 (1974).
- [4] V. I. Iacob *et al.*, *Progress in Research*, Cyclotron Institute, Texas A&M University (2000-2001), p. I-30.